## PANEITZ-TYPE OPERATORS AND APPLICATIONS

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## To the memory of André Lichnerowicz

Given $(M, g)$ a smooth 4-dimensional Riemannian manifold, let $S_{g}$ be the scalar curvature of $g$, and let $R c_{g}$ be the Ricci curvature of $g$. The Paneitz operator, discovered in [21], is the fourth-order operator defined by

$$
P_{g}^{4} u=\Delta_{g}^{2} u-\operatorname{div}_{g}\left(\frac{2}{3} S_{g} g-2 R c_{g}\right) d u
$$

where $\Delta_{g} u=-\operatorname{div}_{g} \nabla u$ is the Laplacian of $u$ with respect to $g$. When $(M, g)$ is the 4-dimensional standard unit sphere $\left(S^{4}, h\right)$, we get that

$$
P_{h}^{4} u=\Delta_{h}^{2} u+2 \Delta_{h} u
$$

The Paneitz operator is conformally invariant in the sense that if $\tilde{g}=e^{2 \varphi} g$ is a conformal metric to $g$, then for all $u \in C^{\infty}(M)$,

$$
P_{\tilde{g}}^{4} u=e^{-4 \varphi} P_{g}^{4}(u)
$$

The 2-dimensional analogue of this relation is

$$
\Delta_{\tilde{g}} u=e^{-2 \varphi} \Delta_{g} u
$$

When the dimension is 2 , it is well known that the scalar curvatures of $g$ and $\tilde{g}$ are related by the equation

$$
\Delta_{g} \varphi+\frac{1}{2} S_{g}=\frac{1}{2} S_{\tilde{g}} e^{2 \varphi}
$$

When the dimension is 4 , we get that

$$
P_{g}^{4} \varphi+Q_{g}^{4}=Q_{\tilde{g}}^{4} e^{4 \varphi}
$$

where

$$
Q_{g}^{4}=\frac{1}{6}\left(\Delta_{g} S_{g}+S_{g}^{2}-3\left|R c_{g}\right|^{2}\right)
$$

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With respect to the Euler-Poincaré characteristic, in dimension 2, we have that

$$
\chi(M)=\frac{1}{4 \pi} \int_{M} S_{g} d v_{g}
$$

while in dimension 4,

$$
\chi(M)=\frac{1}{8 \pi^{2}} \int_{M}\left(\frac{1}{4}\left|W_{g}\right|^{2}+Q_{g}^{4}\right) d v_{g}
$$

where $W_{g}$ stands for the Weyl tensor of $g$. Beautiful works on $P_{g}^{4}$ have been developed recently by Beckner [2], Branson-Chang-Yang [4], Chang-Yang [7], Chang-GurskyYang [6], and Gursky [12]. Also see the survey Chang [5]. The Paneitz operator was generalized to higher dimensions by Branson [3]. Given ( $M, g$ ) a smooth compact Riemannian $n$-manifold, $n \geq 5$, let $P_{g}^{n}$ be the operator defined by

$$
P_{g}^{n} u=\Delta_{g}^{2} u-\operatorname{div}_{g}\left(a_{n} S_{g} g+b_{n} R c_{g}\right) d u+\frac{n-4}{2} Q_{g}^{n} u
$$

where

$$
Q_{g}^{n}=\frac{1}{2(n-1)} \Delta_{g} S_{g}+\frac{n^{3}-4 n^{2}+16 n-16}{8(n-1)^{2}(n-2)^{2}} S_{g}^{2}-\frac{2}{(n-2)^{2}}\left|R c_{g}\right|^{2}
$$

and

$$
\left\{\begin{array}{l}
a_{n}=\frac{(n-2)^{2}+4}{2(n-1)(n-2)} \\
b_{n}=-\frac{4}{n-2}
\end{array}\right.
$$

If $\tilde{g}=\varphi^{4 /(n-4)} g$ is a conformal metric to $g$, then for all $u \in C^{\infty}(M)$,

$$
P_{g}^{n}(u \varphi)=\varphi^{(n+4) /(n-4)} P_{\tilde{g}}^{n} u
$$

In particular,

$$
P_{g}^{n} \varphi=\frac{n-4}{2} Q_{\tilde{g}}^{n} \varphi^{(n+4) /(n-4)}
$$

These two relations have well-known analogues when dealing with the conformal Laplacian $L_{g}^{n}$, the second order operator whose expression is given by

$$
L_{g}^{n} u=\Delta_{g} u+\frac{n-2}{4(n-1)} S_{g} u
$$

If $\tilde{g}=\varphi^{4 /(n-2)} g$ is a conformal metric to $g$, for all $u \in C^{\infty}(M)$, we get that

$$
L_{g}^{n}(u \varphi)=\varphi^{(n+2) /(n-2)} L_{\tilde{g}}^{n}(u)
$$

In particular,

$$
L_{g}^{n} \varphi=\frac{n-2}{4(n-1)} S_{\tilde{g}} \varphi^{(n+2) /(n-2)}
$$

On the standard unit sphere $\left(S^{n}, h\right), n \geq 5$, the expression of $P_{h}^{n}$ is

$$
\begin{equation*}
P_{h}^{n} u=\Delta_{h}^{2} u+c_{n} \Delta_{h} u+d_{n} u \tag{0.1}
\end{equation*}
$$

where

$$
\left\{\begin{array}{l}
c_{n}=\frac{1}{2}\left(n^{2}-2 n-4\right)  \tag{0.2}\\
d_{n}=\frac{n-4}{16} n\left(n^{2}-4\right)
\end{array}\right.
$$

that is, $c_{n}=n(n-1) a_{n}+(n-1) b_{n}, d_{n}=(n-4) Q_{h}^{n} / 2$. Note that $P_{h}^{n}=\left(L_{h}^{n}\right)^{2}-2 L_{h}^{n}$. Given $(M, g)$ a smooth $n$-dimensional compact Riemannian manifold, $n \geq 5$, and $\alpha>0$ real, we let $P_{g}$ be the fourth-order operator defined by

$$
\begin{equation*}
P_{g} u=\Delta_{g}^{2} u+\alpha \Delta_{g} u \tag{0.3}
\end{equation*}
$$

Keeping in mind the expression of $P_{g}^{4}$ on the standard sphere, we refer to $P_{g}$ as a Paneitz-type operator. Results and remarks often shift between operators like $P_{g}$, and Paneitz-Branson-type operators like

$$
\tilde{P}_{g} u=\Delta_{g}^{2} u+\alpha \Delta_{g} u+\beta u
$$

where $\alpha$ and $\beta$ are real numbers. Here we should regard $P_{g}$ as the essential part of $\tilde{P}_{g}$, like the Laplacian is the essential part of the conformal Laplacian, and note that the Paneitz-Branson operator $P_{g}^{n}$ reduces to $\tilde{P}_{g}$ when $g$ is Einstein. A natural space when studying $P_{g}$ is the Sobolev space $H_{2}^{2}(M)$ defined as the completion of $C^{\infty}(M)$ with respect to the norm

$$
\|u\|^{2}=\left\|\nabla^{2} u\right\|_{2}^{2}+\|\nabla u\|_{2}^{2}+\|u\|_{2}^{2}
$$

Following standard notations, $\|\cdot\|_{p}$ in the above expression stands for the $L^{p}$-norm (with respect to the Riemannian measure $d v_{g}$ ). As is well known and easy to see, for all $u \in C^{\infty}(M)$,

$$
\left(\Delta_{g} u\right)^{2} \leq n\left|\nabla^{2} u\right|^{2}
$$

Conversely, by the Bochner-Lichnerowicz-Weitzenböck formula we get that

$$
\begin{aligned}
\int_{M}\left|\nabla^{2} u\right|^{2} d v_{g} & =\int_{M}\left|\Delta_{g} u\right|^{2} d v_{g}-\int_{M} R c_{g}(\nabla u, \nabla u) d v_{g} \\
& \leq \int_{M}\left|\Delta_{g} u\right|^{2} d v_{g}+k \int_{M}|\nabla u|^{2} d v_{g}
\end{aligned}
$$

where $k$ is such that $R c_{g} \geq-k$. Hence, $\|\cdot\|_{H_{2}^{2}}$ defined by

$$
\|u\|_{H_{2}^{2}}^{2}=\int_{M}\left(P_{g} u\right) u d v_{g}+\int_{M} u^{2} d v_{g}
$$

is a norm on $H_{2}^{2}(M)$ that is equivalent to the above more classical one $\|\cdot\|$. Here and in what follows,

$$
\int_{M}\left(P_{g} u\right) u d v_{g}=\int_{M}\left(\Delta_{g} u\right)^{2} d v_{g}+\alpha \int_{M}|\nabla u|^{2} d v_{g}
$$

By the Sobolev embedding theorem (see, e.g., [14] and recall that $n \geq 5$ ), we get an embedding of $H_{2}^{2}(M)$ in $L^{2^{\sharp}}(M)$, where

$$
2^{\sharp}=\frac{2 n}{n-4} .
$$

This embedding is critical. It is also continuous, so that there exist $A \in \mathbb{R}$ and $B \in \mathbb{R}$ such that for all $u \in H_{2}^{2}(M)$,

$$
\begin{equation*}
\|u\|_{2^{\sharp}}^{2} \leq A \int_{M}\left(P_{g} u\right) u d v_{g}+B\|u\|_{2}^{2} . \tag{S1}
\end{equation*}
$$

Another possible inequality is that for all $u \in H_{2}^{2}(M)$,

$$
\begin{equation*}
\|u\|_{2^{\sharp}}^{2} \leq A \int_{M}\left(P_{g} u\right) u d v_{g}+B\|u\|_{H_{1}^{2}}^{2}, \tag{S2}
\end{equation*}
$$

where $\|\cdot\|_{H_{1}^{2}}$ is the usual norm of $H_{1}^{2}(M)$ given by

$$
\|u\|_{H_{1}^{2}}^{2}=\|\nabla u\|_{2}^{2}+\|u\|_{2}^{2}
$$

Given $i=1,2$, we define $A_{\mathrm{opt}}^{(i)}(M)$ as the best constant $A$ in $(S i)$. Formally,

$$
A_{\mathrm{opt}}^{(i)}(M)=\inf \{A \in \mathbb{R} \text { s.t. } \exists B \in \mathbb{R} \text { with the property that }(S i) \text { is valid }\}
$$

where by " $(S i)$ is valid", we mean that (Si) holds with $A$ and $B$ for all $u \in H_{2}^{2}(M)$. Similarly, we define $B_{\mathrm{opt}}^{(i)}(M)$ as the best constant $B$ in $(S i)$, that is,

$$
B_{\mathrm{opt}}^{(i)}(M)=\inf \{B \in \mathbb{R} \text { s.t. } \exists A \in \mathbb{R} \text { with the property that }(S i) \text { is valid }\}
$$

Section 1 of this paper is devoted to the study of these best constants. We determine their exact values and investigate the corresponding optimal inequalities. Applications to a fourth-order partial differential equation of critical growth are presented in Section 2, with special attention devoted to the standard sphere and the Paneitz-Branson
operator $P_{h}^{n}$ given by (0.1). The results we formulate do involve at some point conditions on the parameter $\alpha>0$ in the definition (0.3) of $P_{g}$. It appears that $\alpha$ plays a central role. Results may be expressed equivalently in terms of conditions on $\alpha$ and $\beta$ for the operator

$$
\tilde{P}_{g} u=\Delta_{g}^{2} u+\alpha \Delta_{g} u+\beta u
$$

(and we do so for the operator $P_{h}^{n}$ on the sphere). For the sake of comparison with the classical Sobolev inequalities, and in order to preserve the energy as $\int\left(P_{g} u\right) u d v_{g}$, we deal with definition (0.3) for Paneitz-type operators.

1. Optimal inequalities. As a starting point, let us consider the Euclidean inequality $\forall u \in \mathscr{D}\left(\mathbb{R}^{n}\right), n \geq 5$,

$$
\left(\int_{\mathbb{R}^{n}}|u|^{2^{\sharp}} d x\right)^{2 / 2^{\sharp}} \leq K \int_{\mathbb{R}^{n}}(\Delta u)^{2} d x,
$$

where $\Delta$ stands for the Laplacian with respect to the Euclidean metric. The best constant $K$ in this inequality was studied by Edmunds, Fortunato, and Janelli [10], Lieb [18], and Lions [19]. If $K_{0}$ stands for this best constant, it was shown that

$$
K_{0}^{-1}=\pi^{2} n(n-4)\left(n^{2}-4\right) \Gamma\left(\frac{n}{2}\right)^{4 / n} \Gamma(n)^{-4 / n}
$$

The sharp Euclidean Sobolev inequality is then

$$
\begin{equation*}
\left(\int_{\mathbb{R}^{n}}|u|^{2^{\sharp}} d x\right)^{2 / 2^{\sharp}} \leq K_{0} \int_{\mathbb{R}^{n}}(\Delta u)^{2} d x . \tag{1.1}
\end{equation*}
$$

Its extremal functions are

$$
u_{\epsilon, x_{0}}(x)=\left(\left|x-x_{0}\right|^{2}+\epsilon^{2}\right)^{-(n-4) / 2}
$$

and we have that

$$
\Delta^{2} u_{1, x_{0}}=n(n-4)\left(n^{2}-4\right) u_{1, x_{0}}^{2^{\sharp}} .
$$

Given $\alpha>0$ real, in what follows, we let $P$ be as in ( 0.3 ), with $g=\delta$ being the Euclidean metric. In other words,

$$
P u=\Delta^{2} u+\alpha \Delta u
$$

The first result we prove is the following.
Lemma 1.1. For all $u \in \mathscr{D}\left(\mathbb{R}^{n}\right), n \geq 5$,

$$
\begin{equation*}
\left(\int_{\mathbb{R}^{n}}|u|^{2^{\sharp}} d x\right)^{2 / 2^{\sharp}} \leq K_{0} \int_{\mathbb{R}^{n}}(P u) u d x \tag{1.2}
\end{equation*}
$$

and $K_{0}$ is the best constant in this inequality.

Proof. Suppose that $A>0$ is such that for all $u \in \mathscr{D}\left(\mathbb{R}^{n}\right)$,

$$
\begin{equation*}
\left(\int_{\mathbb{R}^{n}}|u|^{2^{\sharp}} d x\right)^{2 / 2^{\sharp}} \leq A \int_{\mathbb{R}^{n}}(P u) u d x \tag{1.3}
\end{equation*}
$$

Let $\mathscr{B}$ be the unit ball in $\mathbb{R}^{n}$. Then for all $u \in \mathscr{D}(\mathscr{B})$,

$$
\left(\int_{\mathscr{B}}|u|^{2^{\sharp}} d x\right)^{2 / 2^{\sharp}} \leq A \int_{\mathscr{B}}(P u) u d x .
$$

Also let

$$
\begin{aligned}
& H_{0,2}^{2}(\mathscr{B})=\text { completion of } \mathscr{D}(\mathscr{P}) \text { w.r.t. }\|\cdot\|_{H_{2}^{2}}, \\
& H_{0,1}^{2}(\mathscr{B})=\text { completion of } \mathscr{D}(\mathscr{P}) \text { w.r.t. }\|\cdot\|_{H_{1}^{2}},
\end{aligned}
$$

where, as above,

$$
\|u\|_{H_{2}^{2}}^{2}=\int_{\mathscr{B}}(P u) u d x+\int_{\mathscr{B}} u^{2} d x
$$

The embedding of $H_{0,2}^{2}(\mathscr{B})$ in $H_{0,1}^{2}(\mathscr{B})$ is compact, while the embedding of $H_{0,1}^{2}(\mathscr{B})$ in $L^{2}(\mathscr{B})$ is continuous. It easily follows that for any $\epsilon>0$, there exists $B_{\epsilon}>0$ such that for all $u \in \mathscr{D}(\mathscr{B})$,

$$
\|u\|_{H_{1}^{2}}^{2} \leq \epsilon\|u\|_{H_{2}^{2}}^{2}+B_{\epsilon}\|u\|_{2}^{2}
$$

In particular,

$$
\|\nabla u\|_{2}^{2} \leq \frac{\epsilon}{1-\alpha \epsilon}\|\Delta u\|_{2}^{2}+\frac{B_{\epsilon}-1+\epsilon}{1-\alpha \epsilon}\|u\|_{2}^{2}
$$

and we get that for any $\epsilon>0$, there exists $B_{\epsilon}>0$ such that for all $u \in \mathscr{D}(\mathscr{B})$,

$$
\left(\int_{\mathscr{B}}|u|^{2^{\sharp}} d x\right)^{2 / 2^{\sharp}} \leq(A+\epsilon) \int_{\mathscr{B}}(\Delta u)^{2} d x+B_{\epsilon} \int_{\mathscr{B}} u^{2} d x .
$$

Let $\mathscr{B}_{\delta}$ be the ball of center 0 and radius $\delta>0$. For any $\delta<1$, any $\epsilon>0$, and all $u \in \mathscr{D}\left(\mathscr{P}_{\delta}\right)$,

$$
\left(\int_{\mathscr{B}_{\delta}}|u|^{2^{\sharp}} d x\right)^{2 / 2^{\sharp}} \leq(A+\epsilon) \int_{\mathscr{B}_{\delta}}(\Delta u)^{2} d x+B_{\epsilon} \int_{\mathscr{O}_{\delta}} u^{2} d x .
$$

By Hölder's inequality,

$$
\int_{\mathscr{A}_{\delta}} u^{2} d x \leq\left|\mathscr{F}_{\delta}\right|^{1-2 / 2^{\sharp}}\left(\int_{\mathscr{B}_{\delta}}|u|^{2^{\sharp}} d x\right)^{2 / 2^{\sharp}},
$$

where $\left|\mathscr{B}_{\delta}\right|$ stands for the Euclidean volume of $\mathscr{B}_{\delta}$. Choosing $\delta$ small, we are led to the following: for any $\epsilon>0$, there exists $\delta>0$ such that for all $u \in \mathscr{D}\left(\mathscr{B}_{\delta}\right)$,

$$
\left(\int_{\mathscr{B}_{\delta}}|u|^{2^{\sharp}} d x\right)^{2 / 2^{\sharp}} \leq(A+2 \epsilon) \int_{\mathscr{B}_{\delta}}(\Delta u)^{2} d x .
$$

Given $u \in \mathscr{D}\left(\mathbb{R}^{n}\right)$ and $\lambda>0$, let $u_{\lambda}(x)=u(\lambda x)$. For $\lambda$ large, $u_{\lambda} \in \mathscr{D}\left(\mathscr{P}_{\delta}\right)$. In addition,

$$
\left(\int_{\mathscr{B}_{\delta}}\left|u_{\lambda}\right|^{2^{\sharp}} d x\right)^{2 / 2^{\sharp}}=\lambda^{4-n}\left(\int_{\mathscr{B}_{\delta}}|u|^{2^{\sharp}} d x\right)^{2 / 2^{\sharp}},
$$

while

$$
\int_{\mathscr{B}_{\delta}}\left(\Delta u_{\lambda}\right)^{2} d x=\lambda^{4-n} \int_{\mathscr{B}_{\delta}}(\Delta u)^{2} d x .
$$

As a consequence, for all $u \in \mathscr{D}\left(\mathbb{R}^{n}\right)$,

$$
\left(\int_{\mathbb{R}^{n}}|u|^{\nmid^{\sharp}} d x\right)^{2 / 2^{\sharp}} \leq(A+2 \epsilon) \int_{\mathbb{R}^{n}}(\Delta u)^{2} d x,
$$

so that if $A$ is as in (1.3), then $A+2 \epsilon \geq K_{0}$. Since $\epsilon$ is arbitrary, $A \geq K_{0}$. Independently, we clearly get from (1.1) that for all $u \in \mathscr{D}\left(\mathbb{R}^{n}\right)$,

$$
\left(\int_{\mathbb{R}^{n}}|u|^{2^{\sharp}} d x\right)^{2 / 2^{\sharp}} \leq K_{0} \int_{\mathbb{R}^{n}}(P u) u d x,
$$

so that $A=K_{0}$ in (1.3) is an admissible value. Lemma 1.1 is established.
The following lemma easily follows from Lemma 1.1.
Lemma 1.2. Let $(M, g)$ be a smooth, compact, $n$-dimensional Riemannian manifold, $n \geq 5$. Any constant $A$ in (S1) or (S2), whatever the constant B, has to be greater than or equal to $K_{0}$.

Proof. Given $x_{0} \in M$, we consider a geodesic, normal coordinates system at $x_{0}$. We let $B_{x_{0}}(\delta)$ be the ball on which this coordinates system is defined. Given $\epsilon>0$, choosing $\delta$ sufficiently small, and by local comparison of the Riemannian metric with the Euclidean metric in the above coordinates system, we easily get from ( $S 1$ ) and (S2) that for all $u \in \mathscr{D}\left(\mathscr{H}_{\delta}\right)$,

$$
\begin{align*}
& \|u\|_{2^{\sharp}}^{2} \leq(A+\epsilon) \int_{\mathbb{R}^{n}}(P u) u d x+\tilde{B}\|u\|_{2}^{2},  \tag{1.4}\\
& \|u\|_{2^{\sharp}}^{2} \leq(A+\epsilon) \int_{\mathbb{R}^{n}}(P u) u d x+\tilde{B}\|u\|_{H_{1}^{2}}^{2}, \tag{1.5}
\end{align*}
$$

where the norms and $P$ in the above expressions are understood with respect to the Euclidean metric. Passing in the coordinates system, we indeed have that

$$
\left(\Delta_{g} u\right)^{2} \leq(\Delta u)^{2}+\hat{\epsilon}(\delta)\left|\nabla^{2} u\right|^{2}+\hat{\epsilon}(\delta)|\nabla u|^{2},
$$

where $\hat{\epsilon}(\delta) \rightarrow 0$ as $\delta \rightarrow 0$, while by the Bochner-Lichnerowicz-Weitzenböck formula,

$$
\int_{\mathbb{R}^{n}}\left|\nabla^{2} u\right|^{2} d x=\int_{\mathbb{R}^{n}}(\Delta u)^{2} d x
$$

As in the proof of Lemma 1.1, it is enough to deal with (1.4). Applying Hölder's inequality to the $L^{2}$-norm of $u$, with (1.4) as in the proof of Lemma 1.1, we get that for all $\epsilon>0$, there exists $\delta>0$, such that for all $u \in \mathscr{D}\left(\mathscr{B}_{\delta}\right)$,

$$
\|u\|_{2^{\sharp}}^{2} \leq(A+2 \epsilon) \int_{\mathbb{R}^{n}}(P u) u d x .
$$

For homogeneity reasons, see once more the proof of Lemma 1.1; such an inequality, if valid for functions with small support, has to be valid for all functions with compact support. By Lemma 1.1, this implies that $A+2 \epsilon \geq K_{0}$. Since $\epsilon$ is arbitrary, this proves the result.

The answer to the first question we asked, dealing with the exact value of $A_{\mathrm{opt}}^{(i)}(M)$, now follows from Lemma 1.2.

Theorem 1.1. Let $(M, g)$ be a smooth, compact, $n$-dimensional Riemannian manifold, $n \geq 5$, and let $P_{g}$ be as in (0.3), that is,

$$
P_{g} u=\Delta_{g}^{2} u+\alpha \Delta_{g} u
$$

For any $\epsilon>0$, there exists $B_{\epsilon} \in \mathbb{R}$ such that for all $u \in H_{2}^{2}(M)$,

$$
\begin{align*}
& \|u\|_{2^{\sharp}}^{2} \leq(1+\epsilon) K_{0} \int_{M}\left(P_{g} u\right) u d v_{g}+B_{\epsilon}\|u\|_{2}^{2},  \tag{1.6}\\
& \|u\|_{2^{\sharp}}^{2} \leq(1+\epsilon) K_{0} \int_{M}\left(P_{g} u\right) u d v_{g}+B_{\epsilon}\|u\|_{H_{1}^{2}}^{2} . \tag{1.7}
\end{align*}
$$

In particular, $A_{\mathrm{opt}}^{(1)}(M)=A_{\mathrm{opt}}^{(2)}(M)=K_{0}$. Moreover, we can choose $B_{\epsilon}$ depending only on $n, \epsilon, \alpha$, a bound for the Ricci curvature $R c_{g}$ of $g$, and a positive lower bound for the injectivity radius $i_{g}$ of $g$.

Proof. It suffices to prove the result for (1.6), since

$$
\|u\|_{2}^{2} \leq\|u\|_{H_{1}^{2}}^{2}
$$

Let $\Lambda>0$ be such that $\left|R c_{g}\right| \leq \Lambda$, and let $i>0$ be such that $i_{g} \geq i$. By a well-known result of Anderson [1], the $C^{1,1 / 2}$-harmonic radius of $g$ is bounded from below by a positive constant depending mainly on $n, \Lambda$, and $i$. More precisely, given $\delta>0$, there exists $r_{H}=r_{H}(n, \delta, \Lambda, i)$, depending only on these constants, such that the following holds: for any $x \in M$, there exists a harmonic chart $\varphi$ on the ball $B_{x}\left(r_{H}\right)$ with the
property that

$$
\begin{gathered}
\frac{1}{1+\delta} \delta_{i j} \leq g_{i j} \leq(1+\delta) \delta_{i j} \text { as bilinear forms } \\
\sum_{k=1}^{n} r_{H} \sup _{x}\left|\partial_{k} g_{i j}(x)\right|+\sum_{k=1}^{n} r_{H}^{3 / 2} \sup _{y \neq z} \frac{\left|\partial_{k} g_{i j}(y)-\partial_{k} g_{i j}(z)\right|}{d_{g}(y, z)^{1 / 2}} \leq \delta,
\end{gathered}
$$

where $d_{g}$ is the distance with respect to $g$, and the $g_{i j}$ 's stand for the components of $g$ in the harmonic coordinates system. Without loss of generality, we may assume that $r_{H} \leq 1 / 2$. Let $\epsilon>0$ be given. The expression of the Laplacian in harmonic coordinates is

$$
\Delta_{g} u=-g^{i j} \partial_{i j} u
$$

As we easily check, if $\delta=\delta(\epsilon)$ is small enough, for all $x \in M$, and all $u \in$ $\mathscr{D}\left(B_{x}\left(r_{H}\right)\right)$, then

$$
\int_{\mathbb{R}^{n}}\left(P\left(u \circ \varphi^{-1}\right)\right)\left(u \circ \varphi^{-1}\right) d x \leq(1+\epsilon) \int_{M}\left(P_{g} u\right) u d v_{g} .
$$

Similarly, for all $x \in M$, and all $u \in \mathscr{D}\left(B_{x}\left(r_{H}\right)\right)$,

$$
\int_{M}|u|^{2^{\sharp}} d v_{g} \leq(1+\epsilon) \int_{\mathbb{R}^{n}}\left|\left(u \circ \varphi^{-1}\right)\right|^{2^{\sharp}} d x .
$$

It follows from these relations and Lemma 1.1 that for any $\epsilon>0$, there exists $\delta>0$ such that for all $u \in \mathscr{D}\left(B_{x}\left(r_{H}\right)\right)$,

$$
\begin{equation*}
\left(\int_{M}|u|^{2^{\sharp}} d v_{g}\right)^{2 / 2^{\sharp}} \leq\left(1+\frac{\epsilon}{2}\right) K_{0} \int_{M}\left(P_{g} u\right) u d v_{g} . \tag{loc}
\end{equation*}
$$

Since $M$ is compact, it can be covered by a finite number of balls $B_{x_{i}}\left(r_{H} / 2\right), i=$ $1, \ldots, N$. As is by now classical, we may choose these balls such that any point in $M$ has a neighborhood that intersects at most $S_{0}$ of the $B_{x_{i}}\left(r_{H}\right)$ 's, where the integer $S_{0}=S_{0}(n, \epsilon, \Lambda, i)$ depends only on $n, \epsilon, \Lambda$, and $i$. Let $\alpha_{i} \in \mathscr{D}\left(B_{x_{i}}\left(r_{H}\right)\right)$ be such that $0 \leq \alpha_{i} \leq 1$ and $\alpha_{i}=1$ in $B_{x_{i}}\left(r_{H} / 2\right)$. We set

$$
\eta_{i}=\frac{\alpha_{i}^{2}}{\sum_{k} \alpha_{k}^{2}}
$$

Clearly, $\left(\eta_{i}\right)$ is a partition of unity subordinated to the covering $\left(B_{x_{i}}\left(r_{H}\right)\right)$ such that the $\sqrt{\eta_{i}}$ 's are smooth. As we easily check, we may assume that

$$
\left\{\begin{array}{l}
\left|\nabla \sqrt{\eta_{i}}\right| \leq H_{1} \\
\left|\Delta_{g} \sqrt{\eta_{i}}\right| \leq H_{2} \\
\left|\nabla \Delta_{g} \sqrt{\eta_{i}}\right| \leq H_{3}
\end{array}\right.
$$

where $H_{1}, H_{2}$, and $H_{3}$ only depend on $n, \epsilon, \Lambda$, and $i$. Given $u \in C^{\infty}(M)$, we write that

$$
\|u\|_{2^{\sharp}}^{2}=\left\|u^{2}\right\|_{2^{\sharp} / 2} \leq \sum_{i}\left\|\eta_{i} u^{2}\right\|_{2^{\sharp} / 2}=\sum_{i}\left\|\sqrt{\eta_{i}} u\right\|_{2^{\sharp}}^{2} .
$$

With ( $S_{\text {loc }}$ ), it follows that for all $u \in C^{\infty}(M)$,

$$
\begin{aligned}
\left(\int_{M}|u|^{2^{\sharp}} d v_{g}\right)^{2 / 2^{\sharp}} \leq & \left(1+\frac{\epsilon}{2}\right) K_{0} \sum_{i} \int_{M}\left(\Delta_{g}\left(\sqrt{\eta_{i}} u\right)\right)^{2} d v_{g} \\
& +\alpha\left(1+\frac{\epsilon}{2}\right) K_{0} \sum_{i} \int_{M}\left|\nabla\left(\sqrt{\eta_{i}} u\right)\right|^{2} d v_{g} .
\end{aligned}
$$

Hence,

$$
\begin{align*}
\left(\int_{M}|u|^{2^{\sharp}} d v_{g}\right)^{2 / 2^{\sharp}} \leq & \left(1+\frac{\epsilon}{2}\right) K_{0} \sum_{i} \int_{M} \eta_{i}\left(\Delta_{g} u\right)^{2} d v_{g} \\
& +\left(1+\frac{\epsilon}{2}\right) K_{0} \sum_{i} \int_{M}\left(\Delta_{g} \sqrt{\eta_{i}}\right)^{2} u^{2} d v_{g} \\
& +4\left(1+\frac{\epsilon}{2}\right) K_{0} \sum_{i} \int_{M}\left\langle\nabla \sqrt{\eta_{i}},\left.\nabla u\right|_{g} ^{2} d v_{g}\right. \\
& +2\left(1+\frac{\epsilon}{2}\right) K_{0} \sum_{i} \int_{M} \sqrt{\eta_{i}} u\left(\Delta_{g} u\right)\left(\Delta_{g} \sqrt{\eta_{i}}\right) d v_{g} \\
& -4\left(1+\frac{\epsilon}{2}\right) K_{0} \sum_{i} \int_{M} \sqrt{\eta_{i}}\left(\Delta_{g} u\right)\left\langle\nabla \sqrt{\eta_{i}}, \nabla u\right\rangle_{g} d v_{g}  \tag{1.8}\\
& -4\left(1+\frac{\epsilon}{2}\right) K_{0} \sum_{i} \int_{M}\left(\Delta_{g} \sqrt{\eta_{i}}\right) u\left\langle\nabla \sqrt{\eta_{i}}, \nabla u\right\rangle_{g} d v_{g} \\
& +\alpha\left(1+\frac{\epsilon}{2}\right) K_{0} \sum_{i} \int_{M} \eta_{i}|\nabla u|^{2} d v_{g} \\
& +\alpha\left(1+\frac{\epsilon}{2}\right) K_{0} \sum_{i} \int_{M}\left|\nabla \sqrt{\eta_{i}}\right|^{2} u^{2} d v_{g} \\
& +2 \alpha\left(1+\frac{\epsilon}{2}\right) K_{0} \sum_{i} \int_{M} u \sqrt{\eta_{i}}\left|\nabla \sqrt{\eta_{i}}, \nabla u\right\rangle_{g} d v_{g}
\end{align*}
$$

Since $r_{H} \leq 1 / 2$, and if $\delta$ is sufficiently small, we get that for every $x \in M, \varphi\left(B_{x}\left(r_{H}\right)\right)$ $\subset \mathscr{B}$, where $\mathscr{B}$ stands for the unit ball in $\mathbb{R}^{n}$ with center 0 . As in the proof of Lemma 1.1 , given $\epsilon_{0}>0$, there exists $C_{1}$, depending only on $\epsilon_{0}$, such that for all $u \in \mathscr{D}(\mathscr{P})$,

$$
\int_{\mathscr{B}}|\nabla u|^{2} d x \leq \frac{\epsilon_{0}}{2} \int_{\mathscr{B}}(\Delta u)^{2} d x+C_{1} \int_{\mathscr{B}} u^{2} d x
$$

As before, this implies that for all $x \in M$, and all $u \in \mathscr{D}\left(B_{x}\left(r_{H}\right)\right)$,

$$
\int_{M}|\nabla u|^{2} d v_{g} \leq \epsilon_{0} \int_{M}\left(\Delta_{g} u\right)^{2} d v_{g}+C_{2} \int_{M} u^{2} d v_{g}
$$

where $C_{2}$ depends only on $\epsilon_{0}$. Let us now deal with the various terms in (1.8). As a starting point,

$$
\sum_{i} \int_{M} \eta_{i}\left(\Delta_{g} u\right)^{2} d v_{g}=\int_{M}\left(\Delta_{g} u\right)^{2} d v_{g}
$$

and

$$
\sum_{i} \int_{M}\left(\Delta_{g} \sqrt{\eta_{i}}\right)^{2} u^{2} d v_{g} \leq S_{0} H_{2}^{2} \int_{M} u^{2} d v_{g}
$$

To ease the notations, set $\mathscr{B}_{i}=B_{x_{i}}\left(r_{H}\right)$. Then, as is easily checked,

$$
\begin{aligned}
\sum_{i} \int_{M}\left\langle\nabla \sqrt{\eta_{i}},\left.\nabla u\right|_{g} ^{2} d v_{g}\right. & \leq \sum_{i} \int_{M}\left|\nabla \sqrt{\eta_{i}}\right|^{2}|\nabla u|^{2} d v_{g} \\
& \leq H_{1}^{2} \sum_{i} \int_{\mathscr{B}_{i}}|\nabla u|^{2} d v_{g} \\
& \leq H_{1}^{2} \epsilon_{0} \sum_{i} \int_{\mathscr{B}_{i}}\left(\Delta_{g} u\right)^{2} d v_{g}+H_{1}^{2} C_{2} \int_{\mathscr{B}_{i}} u^{2} d v_{g} \\
& \leq S_{0} H_{1}^{2} \epsilon_{0} \int_{M}\left(\Delta_{g} u\right)^{2} d v_{g}+S_{0} H_{1}^{2} C_{2} \int_{M} u^{2} d v_{g}
\end{aligned}
$$

Similarly,

$$
\begin{aligned}
\left|\sum_{i} \int_{M} \sqrt{\eta_{i}} u\left(\Delta_{g} u\right)\left(\Delta_{g} \sqrt{\eta_{i}}\right) d v_{g}\right|= & \mid \sum_{i} \int_{M}\left\langle\nabla \sqrt{\eta_{i}}, \nabla u\right\rangle_{g} u\left(\Delta_{g} \sqrt{\eta_{i}}\right) d v_{g} \\
& +\sum_{i} \int_{M} \sqrt{\eta_{i}}|\nabla u|^{2}\left(\Delta_{g} \sqrt{\eta_{i}}\right)^{2} d v_{g} \\
& +\sum_{i} \int_{M} \sqrt{\eta_{i}} u\left\langle\nabla u, \nabla\left(\Delta_{g} \sqrt{\eta_{i}}\right)\right\rangle_{g} d v_{g} \mid .
\end{aligned}
$$

Hence,

$$
\begin{aligned}
& \left|\sum_{i} \int_{M} \sqrt{\eta_{i}} u\left(\Delta_{g} u\right)\left(\Delta_{g} \sqrt{\eta_{i}}\right) d v_{g}\right| \\
& \quad \leq\left(H_{1} H_{2}+H_{3}\right) \sum_{i} \int_{\mathscr{B}_{i}}|u||\nabla u| d v_{g}+H_{2}^{2} \sum_{i} \int_{\mathscr{B}_{i}}|\nabla u|^{2} d v_{g} \\
& \quad \leq \frac{1}{2}\left(H_{1} H_{2}+H_{3}\right) \sum_{i} \int_{\mathscr{B}_{i}} u^{2} d v_{g}+\left(\frac{1}{2}\left(H_{1} H_{2}+H_{3}\right)+H_{2}^{2}\right) \sum_{i} \int_{\mathscr{B}_{i}}|\nabla u|^{2} d v_{g}
\end{aligned}
$$

As above, this leads to

$$
\begin{aligned}
& \left|\sum_{i} \int_{M} \sqrt{\eta_{i}} u\left(\Delta_{g} u\right)\left(\Delta_{g} \sqrt{\eta_{i}}\right) d v_{g}\right| \\
& \quad \leq\left(\frac{S_{0}}{2}\left(H_{1} H_{2}+H_{3}\right)+S_{0} H_{2}^{2}\right) \epsilon_{0} \int_{M}\left(\Delta_{g} u\right)^{2} d v_{g} \\
& \quad \\
& \quad+\left(\frac{S_{0}}{2}\left(H_{1} H_{2}+H_{3}\right)+\left(\frac{S_{0}}{2}\left(H_{1} H_{2}+H_{3}\right)+S_{0} H_{2}^{2}\right) C_{2}\right) \int_{M} u^{2} d v_{g}
\end{aligned}
$$

On the other hand,

$$
\sum_{i} \int_{M} \sqrt{\eta_{i}}\left(\Delta_{g} u\right)\left\langle\nabla \sqrt{\eta_{i}}, \nabla u\right\rangle_{g} d v_{g}=\frac{1}{2} \sum_{i} \int_{M}\left(\Delta_{g} u\right)\left\langle\nabla \eta_{i}, \nabla u\right\rangle_{g} d v_{g}=0
$$

while

$$
\begin{aligned}
\mid \sum_{i} \int_{M}\left(\Delta_{g} \sqrt{\eta_{i}}\right) & u\left\langle\nabla \sqrt{\eta_{i}}, \nabla u\right\rangle_{g} d v_{g} \mid \\
& \leq H_{1} H_{2} \sum_{i} \int_{\mathscr{B}_{i}}|u||\nabla u| d v_{g} \\
& \leq \frac{H_{1} H_{2}}{2} \sum_{i} \int_{\mathscr{B}_{i}} u^{2} d v_{g}+\frac{H_{1} H_{2}}{2} \sum_{i} \int_{\mathscr{B}_{i}}|\nabla u|^{2} d v_{g}
\end{aligned}
$$

Here again, it follows that

$$
\begin{aligned}
\mid \sum_{i} \int_{M}\left(\Delta_{g} \sqrt{\eta_{i}}\right) u & \left\langle\nabla \sqrt{\eta_{i}}, \nabla u\right\rangle_{g} d v_{g} \mid \\
& \leq \frac{H_{1} H_{2}}{2} S_{0} \epsilon_{0} \int_{M}\left(\Delta_{g} u\right)^{2} d v_{g}+\frac{H_{1} H_{2}}{2} S_{0}\left(1+C_{2}\right) \int_{M} u^{2} d v_{g}
\end{aligned}
$$

Moreover,

$$
\sum_{i} \int_{M} \eta_{i}|\nabla u|^{2} d v_{g} \leq \sum_{i} \int_{\mathscr{B}_{i}}|\nabla u|^{2} d v_{g} \leq S_{0} \epsilon_{0} \int_{M}\left(\Delta_{g} u\right)^{2} d v_{g}+S_{0} C_{2} \int_{M} u^{2} d v_{g}
$$

while

$$
\sum_{i} \int_{M}\left|\nabla \sqrt{\eta_{i}}\right|^{2} u^{2} d v_{g} \leq S_{0} H_{2}^{2} \int_{M} u^{2} d v_{g}
$$

and

$$
\begin{aligned}
& \left|\sum_{i} \int_{M} u \sqrt{\eta_{i}}\left\langle\nabla \sqrt{\eta_{i}}, \nabla u\right\rangle_{g} d v_{g}\right| \\
& \quad \leq H_{2} \sum_{i} \int_{\mathscr{B}_{i}}|u||\nabla u| d v_{g} \leq \frac{H_{2}}{2} \sum_{i} \int_{\mathscr{B}_{i}} u^{2} d v_{g}+\frac{H_{2}}{2} \sum_{i} \int_{\mathscr{B}_{i}}|\nabla u|^{2} d v_{g} \\
& \quad \leq \frac{H_{2}}{2} S_{0} \epsilon_{0} \int_{M}\left(\Delta_{g} u\right)^{2} d v_{g}+\frac{H_{2}}{2} S_{0}\left(1+C_{2}\right) \int_{M} u^{2} d v_{g} .
\end{aligned}
$$

Summarizing our various estimates, we get that for all $u \in C^{\infty}(M)$,

$$
\left(\int_{M}|u|^{2^{\sharp}} d v_{g}\right)^{2 / 2^{\sharp}} \leq\left(\left(1+\frac{\epsilon}{2}\right) K_{0}+C_{3} \epsilon_{0}\right) \int_{M}\left(\Delta_{g} u\right)^{2} d v_{g}+C_{4} \int_{M} u^{2} d v_{g}
$$

where the constant $C_{3}$, explicitly known, depends only on $n, \epsilon, \alpha, \Lambda$, and $i$, and where the constant $C_{4}$, explicitly known, depends only on $n, \epsilon, \alpha, \epsilon_{0}, \Lambda$, and $i$. Choosing $\epsilon_{0}$ such that

$$
C_{3} \epsilon_{0}=\frac{\epsilon}{2} K_{0}
$$

proves the theorem.
By Theorem 1.1, $A_{\mathrm{opt}}^{(1)}(M)=A_{\mathrm{opt}}^{(2)}(M)=K_{0}$. It is natural to ask whether these constants are attained in (S1) and (S2), a question to which we now turn. We do believe that the answer should be positive for what concerns ( $S 2$ ), and a partial result in that direction is given by Theorem 1.2. The answer for $(S 1)$ might be more delicate as shown by Proposition 1.2. As a starting point, we state the following elementary lemma that is used in the proof of Theorem 1.2.

Lemma 1.3. Given $M$ a smooth, compact, $n$-dimensional manifold, $n \geq 5$, and $g_{1}$, $g_{2}=\varphi^{4 /(n-4)} g_{1}$ two conformal metrics on $M$,

$$
\int_{M}\left(P_{g_{2}} u\right) u d v_{g_{2}}=\int_{M}\left(P_{g_{1}}(u \varphi)\right)(u \varphi) d v_{g_{1}}+\mathscr{R}
$$

for all $u \in C^{\infty}(M)$, where $P_{g_{1}}$ and $P_{g_{2}}$ are as in $(0.3)$ with $g=g_{1}, g_{2}$, and where $\mathscr{R}$, explicitly known, is made of lower-order terms in the sense that for all $i=1,2$, and all $u \in C^{\infty}(M)$,

$$
|\mathscr{R}| \leq A_{i} \int_{M}|\nabla u|^{2} d v_{g_{i}}+B_{i} \int_{M} u^{2} d v_{g_{i}}
$$

the constants $A_{i}$ and $B_{i}$ being independent of $u$.
Proof. The proof is simple. We may see the result as an easy consequence of the relation

$$
P_{g_{2}}^{n}(u)=\varphi^{-(n+4) /(n-4)} P_{g_{1}}^{n}(u \varphi) .
$$

We may also proceed by direct computation.

We now prove the following theorem, which yields a partial answer to the question we raised above. As already mentioned, the case where $(M, g)$ is not conformally flat remains open.

Theorem 1.2. Let $(M, g)$ be a smooth, compact, $n$-dimensional Riemannian manifold, $n \geq 5$. If $g$ is conformally flat, then $A_{\mathrm{opt}}^{(2)}(M)=K_{0}$ is attained in (S2). In other words, there exists $B \in \mathbb{R}$ such that

$$
\|u\|_{2^{\sharp}}^{2} \leq K_{0} \int_{M}\left(P_{g} u\right) u d v_{g}+B\|u\|_{H_{1}^{2}}^{2}
$$

for all $u \in H_{2}^{2}(M)$.
Proof. Since $(M, g)$ is compact and conformally flat, it can be covered by a finite number of open subsets $\Omega_{i}, i=1, \ldots, N$, such that
(1) for any $i$, there exists $\varphi_{i} \in C^{\infty}(M), \varphi_{i}>0$, with the property that the metric $g_{i}=\varphi_{i}^{4 /(n-4)} g$ is flat on $\Omega_{i}$;
(2) for any $i$, and if $g_{i}$ is as above, $\left(\Omega_{i}, g_{i}\right)$ is isometric to an open subset of $\mathbb{R}^{n}$ equipped with the Euclidean metric.
Let $\left(\eta_{i}\right)$ be a partition of unity subordinated to the covering $\left(\Omega_{i}\right)$. Without loss of generality, we may assume that for any $i, \sqrt{\eta_{i}} \in C^{\infty}(M)$. Given $u \in C^{\infty}(M)$, write that

$$
\left(\int_{M}|u|^{2^{\sharp}} d v_{g}\right)^{2 / 2^{\sharp}} \leq \sum_{i}\left(\int_{M}\left|\sqrt{\eta_{i}} u\right|^{2^{\sharp}} d v_{g}\right)^{2 / 2^{\sharp}} .
$$

For any $i$, let $\psi_{i}=\varphi_{i}^{-1}$ so that $g=\psi_{i}^{4 /(n-4)} g_{i}$. Then for all $u \in C^{\infty}(M)$,

$$
\int_{M}\left|\sqrt{\eta_{i}} u\right|^{2^{\sharp}} d v_{g}=\int_{M}\left|\sqrt{\eta_{i}} u \psi_{i}\right|^{2^{\sharp}} d v_{g_{i}} .
$$

It follows from Lemma 1.1 that

$$
\left(\int_{M}\left|\sqrt{\eta_{i}} u \psi_{i}\right|^{2^{\sharp}} d v_{g_{i}}\right)^{2 / 2^{\sharp}} \leq K_{0} \int_{M}\left(P_{g_{i}}\left(\sqrt{\eta_{i}} u \psi_{i}\right)\right)\left(\sqrt{\eta_{i}} u \psi_{i}\right) d v_{g_{i}} .
$$

By Lemma 1.3, for any $i$, and for all $u \in C^{\infty}(M)$,

$$
\begin{aligned}
\left(\int_{M}\left|\sqrt{\eta_{i}} u\right|^{2^{\sharp}} d v_{g}\right)^{2 / 2^{\sharp}} \leq & K_{0} \int_{M} P_{g}\left(\sqrt{\eta_{i}} u\right)\left(\sqrt{\eta_{i}} u\right) d v_{g} \\
& +A_{i} \int_{M}|\nabla u|^{2} d v_{g}+B_{i} \int_{M} u^{2} d v_{g}
\end{aligned}
$$

Hence, for all $u \in C^{\infty}(M)$,

$$
\begin{aligned}
\left(\int_{M}|u|^{2^{\sharp}} d v_{g}\right)^{2 / 2^{\sharp}} \leq & K_{0} \sum_{i} \int_{M} P_{g}\left(\sqrt{\eta_{i}} u\right)\left(\sqrt{\eta_{i}} u\right) d v_{g} \\
& +\left(\sum_{i} A_{i}\right) \int_{M}|\nabla u|^{2} d v_{g}+\left(\sum_{i} B_{i}\right) \int_{M} u^{2} d v_{g}
\end{aligned}
$$

Now, as in the proof of Theorem 1.1,

$$
\begin{aligned}
& \sum_{i} \int_{M} P_{g}\left(\sqrt{\eta_{i}} u\right)\left(\sqrt{\eta_{i}} u\right) d v_{g} \\
& \leq \sum_{i} \int_{M}\left(\Delta_{g}\left(\sqrt{\eta_{i}} u\right)\right)^{2} d v_{g}+\alpha \sum_{i} \int_{M}\left|\nabla\left(\sqrt{\eta_{i}} u\right)\right|^{2} d v_{g} \\
& \leq \sum_{i} \int_{M} \eta_{i}\left(\Delta_{g} u\right)^{2} d v_{g}+\sum_{i} \int_{M}\left(\Delta_{g} \sqrt{\eta_{i}}\right)^{2} u^{2} d v_{g} \\
&+4 \sum_{i} \int_{M}\left\langle\nabla \sqrt{\eta_{i}},\left.\nabla u\right|_{g} ^{2} d v_{g}+2 \sum_{i} \int_{M} \sqrt{\eta_{i}} u\left(\Delta_{g} u\right)\left(\Delta_{g} \sqrt{\eta_{i}}\right) d v_{g}\right. \\
&-4 \sum_{i} \int_{M} \sqrt{\eta_{i}}\left(\Delta_{g} u\right)\left\langle\nabla \sqrt{\eta_{i}}, \nabla u\right\rangle_{g} d v_{g} \\
&-4 \sum_{i} \int_{M}\left(\Delta_{g} \sqrt{\eta_{i}}\right) u\left\langle\nabla \sqrt{\eta_{i}}, \nabla u\right\rangle_{g} d v_{g} \\
&+\alpha \sum_{i} \int_{M} \eta_{i}|\nabla u|^{2} d v_{g}+\alpha \sum_{i} \int_{M}\left|\nabla \sqrt{\eta_{i}}\right|^{2} u^{2} d v_{g} \\
&+2 \alpha \sum_{i} \int_{M} u \sqrt{\eta_{i}}\left\langle\nabla \sqrt{\eta_{i}}, \nabla u\right\rangle_{g} d v_{g}
\end{aligned}
$$

We again analyze the various terms in this expression. First,

$$
\sum_{i} \int_{M} \eta_{i}\left(\Delta_{g} u\right)^{2} d v_{g}=\int_{M}\left(\Delta_{g} u\right)^{2} d v_{g}
$$

while

$$
\sum_{i} \int_{M}\left(\Delta_{g} \sqrt{\eta_{i}}\right)^{2} u^{2} d v_{g} \leq N\left(\max _{i} \max _{M}\left|\Delta_{g} \sqrt{\eta_{i}}\right|\right)^{2} \int_{M} u^{2} d v_{g}
$$

Similarly,

$$
\left.\sum_{i} \int_{M}\left\langle\nabla \sqrt{\eta_{i}},\left.\nabla u\right|_{g} ^{2} d v_{g} \leq N\left(\max _{i} \max _{M}\left|\nabla \sqrt{\eta_{i}}\right|\right)^{2} \int_{M}\right| \nabla u\right|^{2} d v_{g}
$$

On the other hand,

$$
\begin{aligned}
\sum_{i} \int_{M} \sqrt{\eta_{i}} u\left(\Delta_{g} u\right)\left(\Delta_{g} \sqrt{\eta_{i}}\right) d v_{g}= & \sum_{i} \int_{M}\left\langle\nabla \sqrt{\eta_{i}}, \nabla u\right\rangle_{g} u\left(\Delta_{g} \sqrt{\eta_{i}}\right) d v_{g} \\
& +\sum_{i} \int_{M} \sqrt{\eta_{i}}|\nabla u|^{2}\left(\Delta_{g} \sqrt{\eta_{i}}\right)^{2} d v_{g} \\
& +\sum_{i} \int_{M} \sqrt{\eta_{i}} u\left\langle\nabla u, \nabla\left(\Delta_{g} \sqrt{\eta_{i}}\right)\right\rangle_{g} d v_{g}
\end{aligned}
$$

so that

$$
\begin{aligned}
&\left|\sum_{i} \int_{M} \sqrt{\eta_{i}} u\left(\Delta_{g} u\right)\left(\Delta_{g} \sqrt{\eta_{i}}\right) d v_{g}\right| \\
& \leq \frac{N}{2}\left(\max _{i} \max _{M}\left|\nabla \sqrt{\eta_{i}}\right|\right)\left(\max _{i} \max _{M}\left|\Delta_{g} \sqrt{\eta_{i}}\right|\right) \int_{M} u^{2} d v_{g} \\
&+\frac{N}{2}\left(\max _{i} \max _{M}\left|\nabla \sqrt{\eta_{i}}\right|\right)\left(\max _{i} \max _{M}\left|\Delta_{g} \sqrt{\eta_{i}}\right|\right) \int_{M}|\nabla u|^{2} d v_{g} \\
&+N\left(\max _{i} \max _{M}\left|\Delta_{g} \sqrt{\eta_{i}}\right|\right)^{2} \int_{M}|\nabla u|^{2} d v_{g} \\
&+\frac{N}{2}\left(\max _{i} \max _{M}\left|\nabla \Delta_{g} \sqrt{\eta_{i}}\right|\right) \int_{M} u^{2} d v_{g} \\
&+\frac{N}{2}\left(\max _{i} \max _{M}\left|\nabla \Delta_{g} \sqrt{\eta_{i}}\right|\right) \int_{M}|\nabla u|^{2} d v_{g}
\end{aligned}
$$

Furthermore,

$$
\sum_{i} \int_{M} \sqrt{\eta_{i}}\left(\Delta_{g} u\right)\left\langle\nabla \sqrt{\eta_{i}}, \nabla u\right\rangle_{g} d v_{g}=\frac{1}{2} \sum_{i} \int_{M}\left(\Delta_{g} u\right)\left\langle\nabla \eta_{i}, \nabla u\right\rangle_{g} d v_{g}=0
$$

while

$$
\begin{aligned}
\mid \sum_{i} & \int_{M}\left(\Delta_{g} \sqrt{\eta_{i}}\right) u\left\langle\nabla \sqrt{\eta_{i}}, \nabla u\right\rangle_{g} d v_{g} \mid \\
\leq & \frac{N}{2}\left(\max _{i} \max _{M}\left|\Delta_{g} \sqrt{\eta_{i}}\right|\right)\left(\max _{i} \max _{M}\left|\nabla \sqrt{\eta_{i}}\right|\right) \int_{M} u^{2} d v_{g} \\
& +\frac{N}{2}\left(\max _{i} \max _{M}\left|\Delta_{g} \sqrt{\eta_{i}}\right|\right)\left(\max _{i} \max _{M}\left|\nabla \sqrt{\eta_{i}}\right|\right) \int_{M}|\nabla u|^{2} d v_{g}
\end{aligned}
$$

For the next terms, note that

$$
\sum_{i} \int_{M} \eta_{i}|\nabla u|^{2} d v_{g}=\int_{M}|\nabla u|^{2} d v_{g}
$$

and

$$
\sum_{i} \int_{M}\left|\nabla \sqrt{\eta_{i}}\right|^{2} u^{2} d v_{g} \leq N\left(\max _{i} \max _{M}\left|\nabla \sqrt{\eta_{i}}\right|\right)^{2} \int_{M} u^{2} d v_{g}
$$

Finally,

$$
\begin{aligned}
\left.\left|\sum_{i} \int_{M} u \sqrt{\eta_{i}}\right| \nabla \sqrt{\eta_{i}}, \nabla u\right\rangle_{g} d v_{g} \mid \leq & \frac{N}{2}\left(\max _{i} \max _{M}\left|\nabla \sqrt{\eta_{i}}\right|\right) \int_{M} u^{2} d v_{g} \\
& +\frac{N}{2}\left(\max _{i} \max _{M}\left|\nabla \sqrt{\eta_{i}}\right|\right) \int_{M}|\nabla u|^{2} d v_{g}
\end{aligned}
$$

As a conclusion, for all $u \in C^{\infty}(M)$,

$$
\left(\int_{M}|u|^{2^{\sharp}} d v_{g}\right)^{2 / 2^{\sharp}} \leq K_{0} \int_{M}\left(\Delta_{g} u\right)^{2} d v_{g}+A \int_{M}|\nabla u|^{2} d v_{g}+B \int_{M} u^{2} d v_{g},
$$

where $A$ and $B$ are constants that do not depend on $u$. In particular, and for all $u \in C^{\infty}(M)$,

$$
\|u\|_{2^{\sharp}}^{2} \leq K_{0} \int_{M}\left(P_{g} u\right) u d v_{g}+C\|u\|_{H_{1}^{2}}^{2},
$$

where $C$ does not depend on $u$. Theorem 1.2 is thus proved.
When dealing with conformally flat manifolds, Theorem 1.2 indicates that $A_{\mathrm{opt}}^{(2)}(M)$ $=K_{0}$ is attained in (S2). The situation for (S1) is more complicated. In particular, $\alpha$ in the definition (0.3) of $P_{g}$ has to play a role. This is what we prove in Proposition 1.2 below. As a first result, we establish the following.

Proposition 1.1. Let $\left(S^{n}, h\right)$ be the standard unit sphere of $\mathbb{R}^{n+1}, n \geq 5$, and let $P_{h}^{n}$ be the Paneitz-Branson operator given by (0.1), that is,

$$
P_{h}^{n} u=\Delta_{h}^{2} u+c_{n} \Delta_{h} u+d_{n} u
$$

Then

$$
\begin{equation*}
\inf _{u \in C^{\infty}\left(S^{n}\right) \backslash\{0\}} \frac{\int_{S^{n}}\left(P_{h}^{n} u\right) u d v_{h}}{\left(\int_{S^{n}}|u|^{2^{\sharp}} d v_{h}\right)^{2 / 2^{\sharp}}}=\frac{1}{K_{0}} \tag{1.9}
\end{equation*}
$$

with the additional property that for any $\beta>1$, and any $x_{0} \in S^{n}$, if

$$
u_{\beta}=\frac{\left(\beta^{2}-1\right)^{(n-4) / 4}}{(\beta-\cos r)^{(n-4) / 2}}
$$

where $r$ is the distance on $S^{n}$ to $x_{0}$, then

$$
P_{h}^{n} u_{\beta}=d_{n} u_{\beta}^{2^{\sharp}-1}
$$

and $u_{\beta}$ realizes the infimum in the left-hand side of (1.9).

Proof. We first prove (1.9). For that purpose, let $x_{0}$ be some point on $S^{n}$, and let $\Phi: S^{n} \backslash\left\{x_{0}\right\} \rightarrow \mathbb{R}^{n}$ be the stereographic projection of pole $x_{0}$. If $\delta$ stands for the Euclidean metric of $\mathbb{R}^{n}$, then

$$
\left(\Phi^{-1}\right)^{\star} h=\varphi^{4 /(n-4)} \delta,
$$

where

$$
\varphi(x)=\left(\frac{4}{\left(1+|x|^{2}\right)^{2}}\right)^{(n-4) / 4}
$$

By conformal invariance of the Paneitz-Branson operator $P_{g}^{n}$, we get that for all $u \in \mathscr{D}\left(\mathbb{R}^{n}\right)$,

$$
\begin{equation*}
\frac{\int_{\mathbb{R}^{n}}\left(\left(\Delta_{\tilde{h}} u\right)^{2}+c_{n}|\nabla u|^{2}+d_{n} u^{2}\right) d v_{\tilde{h}}}{\left(\int_{\mathbb{R}^{n}}|u|^{2^{\sharp}} d v_{\tilde{h}}\right)^{2 / 2^{\sharp}}}=\frac{\int_{\mathbb{R}^{n}}(\Delta(u \varphi))^{2} d x}{\left(\int_{\mathbb{R}^{n}}|u \varphi|^{2^{\sharp}} d x\right)^{2 / 2^{\sharp}}}, \tag{1.10}
\end{equation*}
$$

where $\tilde{h}=\left(\Phi^{-1}\right)^{\star} h$. Suppose now that

$$
\begin{equation*}
\inf _{u \in C^{\infty}\left(S^{n}\right) \backslash\{0\}} \frac{\int_{S^{n}}\left(\left(\Delta_{h} u\right)^{2}+c_{n}|\nabla u|^{2}+d_{n} u^{2}\right) d v_{h}}{\left(\int_{S^{n}}|u|^{2^{\sharp}} d v_{h}\right)^{2 / 2^{\sharp}}}<\frac{1}{K_{0}}, \tag{1.11}
\end{equation*}
$$

and let $u_{0} \in C^{\infty}\left(S^{n}\right), u_{0} \not \equiv 0$, be such that

$$
\frac{\int_{S^{n}}\left(\left(\Delta_{h} u_{0}\right)^{2}+c_{n}\left|\nabla u_{0}\right|^{2}+d_{n} u_{0}^{2}\right) d v_{h}}{\left(\int_{S^{n}}\left|u_{0}\right|^{2^{\sharp}} d v_{h}\right)^{2 / 2^{\sharp}}}<\frac{1}{K_{0}} .
$$

We let $\left(\eta_{s}\right), s>0$ small, be a family of smooth functions on $S^{n}$ having the property that $0 \leq \eta_{s} \leq 1, \eta_{s}=0$ on $B_{P}(s), \eta_{s}=1$ on $S^{n} \backslash B_{P}(2 s)$, and

$$
\left\{\begin{array}{l}
\left|\nabla \eta_{s}\right| \leq \frac{C_{1}}{s} \\
\left|\Delta_{h} \eta_{s}\right| \leq \frac{C_{2}}{s^{2}}
\end{array}\right.
$$

where $C_{1}, C_{2}$ are positive constants that do not depend on $s$. In order to get such a family, we might fix some $\eta_{s_{0}}$ as above, for instance, radially symmetric, and set then, for $s \leq s_{0}, \eta_{s}=\eta_{s_{0}}(r / s)$. As we easily check,

$$
\lim _{s \rightarrow 0} \frac{\int_{S^{n}}\left(\left(\Delta_{h} u_{s}\right)^{2}+c_{n}\left|\nabla u_{s}\right|^{2}+d_{n} u_{s}^{2}\right) d v_{h}}{\left(\int_{S^{n}}\left|u_{S}\right|^{2^{\sharp}} d v_{h}\right)^{2 / 2^{\sharp}}}=\frac{\int_{S^{n}}\left(\left(\Delta_{h} u_{0}\right)^{2}+c_{n}\left|\nabla u_{0}\right|^{2}+d_{n} u_{0}^{2}\right) d v_{h}}{\left(\int_{S^{n}}\left|u_{0}\right|^{2^{\sharp}} d v_{h}\right)^{2 / 2^{\sharp}}},
$$

where $u_{s}=\eta_{s} u_{0}$. Just note here that

$$
\lim _{s \rightarrow 0} \frac{1}{s^{2}} V_{h}\left(B_{P}(2 s) \backslash B_{P}(s)\right)=0,
$$

where $V_{h}(\Omega)$ stands for the volume of $\Omega$ with respect to $h$. Choosing $s$ sufficiently small, it follows from (1.10) and (1.11) that there exists $\tilde{u}_{s} \in \mathscr{D}\left(\mathbb{R}^{n}\right)$ of the form

$$
\tilde{u}_{s}=\left(u_{s} \circ \Phi^{-1}\right) \varphi
$$

such that

$$
\frac{\int_{\mathbb{R}^{n}}\left(\Delta \tilde{u}_{s}\right)^{2} d x}{\left(\int_{\mathbb{R}^{n}}\left|\tilde{u}_{s}\right|^{2^{\sharp}} d x\right)^{2 / 2^{\sharp}}}<\frac{1}{K_{0}} .
$$

This contradicts (1.1) so that (1.9) follows. Now let $u_{\beta}$ be as in the statement of the proposition. As is well known, see, for instance, [15], there exists a conformal diffeomorphism $\varphi_{\beta}$ of $\left(S^{n}, h\right)$ such that

$$
\varphi_{\beta}^{\star} h=u_{\beta}^{4 /(n-4)} h
$$

By conformal invariance of the Paneitz-Branson operator $P_{g}^{n}$, this implies that

$$
P_{h}^{n} u_{\beta}=d_{n} u_{\beta}^{2^{\sharp}-1}
$$

On the other hand, it is easily seen that

$$
d_{n} K_{0}=\omega_{n}^{-4 / n}
$$

where $\omega_{n}$ stands for the volume of $S^{n}$ with respect to $h$. This follows from the relations

$$
4^{n / 2} \omega_{n-1}=\frac{2 \Gamma(n)}{\omega_{n-1} \Gamma(n / 2)^{2}}
$$

and

$$
\omega_{n-1} \Gamma\left(\frac{n}{2}\right)=2 \pi^{n / 2}
$$

Noting that

$$
\int_{S^{n}} u_{\beta}^{2^{\sharp}} d v_{h}=\omega_{n},
$$

we get that

$$
\int_{S^{n}}\left(P_{h}^{n} u_{\beta}\right) u_{\beta} d v_{h}=\frac{1}{K_{0}}\left(\int_{S^{n}} u_{\beta}^{2^{\sharp}} d v_{h}\right)^{2 / 2^{\sharp}} .
$$

This ends the proof of the proposition.
As an ending result in the study of the best first constant, we now prove the following. Since ( $S^{n}, h$ ) is conformally flat, this result has to be compared to Theorem 1.2. Due to the lack of concentration, the approach we use does not allow us to conclude when $n=5$.

Proposition 1.2. Let $\left(S^{n}, h\right)$ be the standard unit sphere of $\mathbb{R}^{n+1}, n \geq 6$, and let $P_{h}$ be as in (0.3) with $g=h$, that is,

$$
P_{h} u=\Delta_{h}^{2} u+\alpha \Delta_{h} u
$$

There exists $B \in \mathbb{R}$ such that for all $u \in H_{2}^{2}\left(S^{n}\right)$,

$$
\|u\|_{2^{\sharp}}^{2} \leq K_{0} \int_{S^{n}}\left(P_{h} u\right) u d v_{h}+B\|u\|_{2}^{2}
$$

if and only if $\alpha \geq c_{n}$, where $c_{n}$ is as in (0.2).
Proof. If $\alpha \geq c_{n}$, the result follows from (1.9), and we may take $B=K_{0} d_{n}$. Suppose, on the contrary, that $\alpha<c_{n}$. For $\beta>1$ real, and $r$ the distance on $S^{n}$ to a given point, we let $u_{\beta}$ be as in Proposition 1.1, that is,

$$
u_{\beta}=\frac{\left(\beta^{2}-1\right)^{(n-4) / 4}}{(\beta-\cos r)^{(n-4) / 2}}
$$

Then, for any $\beta>1$,

$$
\frac{\int_{S^{n}}\left(P_{h}^{n} u_{\beta}\right) u_{\beta} d v_{h}}{\left(\int_{S^{n}} u_{\beta}^{2^{\sharp}} d v_{h}\right)^{2 / 2^{\sharp}}}=\frac{1}{K_{0}} .
$$

Let $B$ be given. Writing that

$$
P_{h} u+B u=P_{h}^{n} u+\left(\alpha-c_{n}\right) \Delta_{h} u+\left(B-d_{n}\right) u
$$

and since

$$
\int_{S^{n}} u_{\beta}^{2^{\sharp}} d v_{h}=\omega_{n},
$$

it follows that for any $\beta>1$,

$$
\begin{align*}
& \frac{\int_{S^{n}}\left(P_{h} u_{\beta}\right) u_{\beta} d v_{h}+B\left\|u_{\beta}\right\|_{2}^{2}}{\left\|u_{\beta}\right\|_{2^{\sharp}}^{2}}  \tag{1.12}\\
& \quad=\frac{1}{K_{0}}+\frac{1}{\omega_{n}^{2 / 2^{\sharp}}}\left(\left(\alpha-c_{n}\right) \int_{S^{n}}\left|\nabla u_{\beta}\right|^{2} d v_{h}+\left(B-d_{n}\right) \int_{S^{n}} u_{\beta}^{2} d v_{h}\right) .
\end{align*}
$$

Performing the changes of variables $x=\tan (r / 2)$ and $y=\sqrt{(\beta+1) /(\beta-1)} x$, it is easily seen that

$$
\int_{S^{n}}\left|\nabla u_{\beta}\right|^{2} d v_{h}=C_{1}(\beta, n)(\beta-1) \int_{0}^{+\infty} \frac{y^{n+1} d y}{\left(1+((\beta-1) /(\beta+1)) y^{2}\right)^{4}\left(1+y^{2}\right)^{n-2}}
$$

and that

$$
\begin{aligned}
\int_{S^{n}} u_{\beta}^{2} d v_{h}= & C_{2}(\beta, n)(\beta-1) \\
& \times \int_{0}^{+\infty} \frac{y^{n-1} d y}{\left(1+((\beta-1) /(\beta+1)) y^{2}\right)^{3}\left(1 /(\beta-1)+(1 /(\beta+1)) y^{2}\right)\left(1+y^{2}\right)^{n-4}},
\end{aligned}
$$

where

$$
C_{1}(\beta, n)=\frac{2^{n}(n-4)^{2} \omega_{n-1}}{(\beta+1)^{3}}, \quad C_{2}(\beta, n)=\frac{2^{n} \omega_{n-1}}{(\beta+1)^{2}}
$$

By the dominated convergence theorem, if $n>6$,

$$
\lim _{\beta \rightarrow 1^{+}} \int_{0}^{+\infty} \frac{y^{n-1} d y}{\left(1+((\beta-1) /(\beta+1)) y^{2}\right)^{3}\left(1 /(\beta-1)+(1 /(\beta+1)) y^{2}\right)\left(1+y^{2}\right)^{n-4}}=0
$$

while

$$
\lim _{\beta \rightarrow 1^{+}} \int_{0}^{+\infty} \frac{y^{n+1} d y}{\left(1+((\beta-1) /(\beta+1)) y^{2}\right)^{4}\left(1+y^{2}\right)^{n-2}}=\int_{0}^{+\infty} \frac{y^{n+1} d y}{\left(1+y^{2}\right)^{n-2}}
$$

The latter integral is a finite positive constant. It follows that if $\alpha<c_{n}$ and $n>6$, then for $\beta>1$ sufficiently close to 1 ,

$$
\left(\alpha-c_{n}\right) \int_{S^{n}}\left|\nabla u_{\beta}\right|^{2} d v_{h}+\left(B-d_{n}\right) \int_{S^{n}} u_{\beta}^{2} d v_{h}<0
$$

We then get with (1.12) that for $\beta>1$ sufficiently close to 1 ,

$$
\int_{S^{n}}\left(P_{h} u_{\beta}\right) u_{\beta} d v_{h}+B\left\|u_{\beta}\right\|_{2}^{2}<\frac{1}{K_{0}}\left\|u_{\beta}\right\|_{2^{\sharp}}^{2} .
$$

This proves the proposition when $n>6$. If $n=6$, we decompose our integrals into three pieces by writing that $\int_{0}^{+\infty}=\int_{0}^{1}+\int_{1}^{1 / \sqrt{\beta-1}}+\int_{1 / \sqrt{\beta-1}}^{+\infty}$. Easy computations then give us that

$$
A_{1}(\beta-1) \ln \left(\frac{1}{\beta-1}\right) \leq \int_{S^{6}}\left|\nabla u_{\beta}\right|^{2} d v_{h} \leq A_{2}(\beta-1) \ln \left(\frac{1}{\beta-1}\right)
$$

for some positive constants $A_{1}<A_{2}$ independent of $\beta$, while

$$
\lim _{\beta \rightarrow 1^{+}} \frac{\int_{S^{6}} u_{\beta}^{2} d v_{h}}{(\beta-1) \ln (1 /(\beta-1))}=0
$$

As above, this gives that if $\alpha<c_{6}$, then for $\beta>1$ sufficiently close to 1 ,

$$
\left(\alpha-c_{6}\right) \int_{S^{6}}\left|\nabla u_{\beta}\right|^{2} d v_{h}+\left(B-d_{6}\right) \int_{S^{6}} u_{\beta}^{2} d v_{h}<0
$$

The proposition is proved.

Parallel with the study of the best first constant, we can ask similar questions on the best second constant. We state our results on this parallel program without any proofs. Details on these proofs can be found in Djadli-Hebey-Ledoux [9]. As a starting point, it is easily seen that whatever $(M, g)$ smooth and compact of dimension $n \geq 5$ is,

$$
B_{\mathrm{opt}}^{(1)}(M)=B_{\mathrm{opt}}^{(2)}(M)=V_{g}^{-4 / n},
$$

where $V_{g}$ is the volume of $(M, g)$. Moreover, we can prove that these constants are attained in the sense that there always exists $A \in \mathbb{R}$ such that for all $u \in H_{2}^{2}(M)$,

$$
\begin{equation*}
\|u\|_{2^{\sharp}}^{2} \leq A \int_{M}\left(P_{g} u\right) u d v_{g}+V_{g}^{-4 / n}\|u\|_{2}^{2}, \tag{1.13}
\end{equation*}
$$

and such that for all $u \in H_{2}^{2}(M)$,

$$
\|u\|_{2^{\sharp}}^{2} \leq A \int_{M}\left(P_{g} u\right) u d v_{g}+V_{g}^{-4 / n}\|u\|_{H_{1}^{2}}^{2} .
$$

There $A$ can be chosen such that it depends only on $n, \alpha$, a lower bound on the Ricci curvature of $g$, a lower bound on the volume of $M$ with respect to $g$, and an upper bound on the diameter of $M$ with respect to $g$. Looking for more precise information on the remaining constant $A$, an easy statement is that $A$ in (1.13) has to be such that

$$
\begin{equation*}
A \geq \frac{2^{\sharp}-2}{\lambda_{1}\left(\lambda_{1}+\alpha\right)} V_{g}^{-4 / n} \tag{1.14}
\end{equation*}
$$

where $\lambda_{1}$ is the first nonzero eigenvalue of $\Delta_{g}$. In the specific case of the standard unit sphere $\left(S^{n}, h\right)$, as proved by Beckner [2], the Sobolev inequality

$$
\|u\|_{p}^{2} \leq \frac{p-2}{n} \omega_{n}^{2 / p-1}\|\nabla u\|_{2}^{2}+\omega_{n}^{2 / p-1}\|u\|_{2}^{2}
$$

holds for all $p \in\left[2,2^{\star}\right]$. By the variational characterization of the first nonzero eigenvalue $\lambda_{1}$ of $\Delta_{h}$, and the Bochner-Lichnerowicz-Weitzenböck formula, it follows that for every $u \in H_{2}^{2}\left(S^{n}\right)$,

$$
\|u\|_{p}^{2} \leq \frac{p-2}{n(n+\alpha)} \omega_{n}^{2 / p-1} \int_{S^{n}}\left(P_{h} u\right) u d v_{h}+\omega_{n}^{2 / p-1}\|u\|_{2}^{2}
$$

It is natural to question whether or not this Beckner's type inequality extends to real numbers $p$ such that $p>2^{\star}$. Assuming that this is the case, and, in particular, that the inequality holds for all $p \in\left[2,2^{\sharp}\right]$, we would get that

$$
\begin{equation*}
\|u\|_{2^{\sharp}}^{2} \leq \frac{8}{n(n-4)(n+\alpha)} \omega_{n}^{-4 / n} \int_{S^{n}}\left(P_{h} u\right) u d v_{h}+\omega_{n}^{-4 / n}\|u\|_{2}^{2} . \tag{1.15}
\end{equation*}
$$

Observe that the first constant in this inequality is the constant given by (1.14) when the manifold considered is the standard sphere. Let $c_{n}$ be as in ( 0.2 ). We can prove (see [9] for details) that if $\alpha \leq c_{n}$, then for all $u \in H_{2}^{2}\left(S^{n}\right)$,

$$
\|u\|_{2^{\sharp}}^{2} \leq \frac{8}{n(n-4)(n+\alpha)} \omega_{n}^{-4 / n} \int_{S^{n}}\left(P_{h} u\right) u d v_{h}+\omega_{n}^{-4 / n}\|u\|_{2}^{2},
$$

and the two constants in this inequality cannot be lowered. In a similar way, we can prove that if $\alpha>c_{n}$, then for all $u \in H_{2}^{2}\left(S^{n}\right)$,

$$
\|u\|_{2^{\sharp}}^{2} \leq \frac{16}{n(n-4)\left(n^{2}-4\right)} \omega_{n}^{-4 / n} \int_{S^{n}}\left(P_{h} u\right) u d v_{h}+\omega_{n}^{-4 / n}\|u\|_{2}^{2},
$$

and again the two constants in this inequality cannot be lowered. In particular, (1.15) is true if $\alpha \leq c_{n}$, but false if $\alpha>c_{n}$. It follows that Beckner's inequality does not extend to $p=2^{\sharp}$. As an ending remark, coming back to an arbitrary, smooth, compact, Riemannian manifold of dimension $n \geq 5$, we mention that it is possible to prove that if $g$ is such that $R c_{g} \geq n-1$, then for all $u \in H_{2}^{2}(M)$,

$$
\begin{equation*}
\|u\|_{2^{\sharp}}^{2} \leq A V_{g}^{-4 / n} \int_{M}\left(P_{g} u\right) u d v_{g}+V_{g}^{-4 / n}\|u\|_{2}^{2}, \tag{1.16}
\end{equation*}
$$

where $A=A(n, \alpha)$, explicitly known, depends only on $n$ and $\alpha$. Let $\hat{A}(n, \alpha)$ be the constant involved in the above inequalities on the sphere

$$
\hat{A}(n, \alpha)= \begin{cases}\frac{8}{n(n-4)(n+\alpha)} & \text { if } \alpha \leq c_{n} \\ \frac{16}{n(n-4)\left(n^{2}-4\right)} & \text { if } \alpha>c_{n}\end{cases}
$$

With respect to what was proved by Ilias [16] when dealing with the Sobolev space $H_{1}^{2}$, an open question we are left with is whether or not (1.16) holds with $A=\hat{A}(n, \alpha)$ when $g$ is such that $R c_{g} \geq n-1$.
2. On a fourth-order partial differential equation. Let $(M, g)$ be a smooth, compact, $n$-dimensional Riemannian manifold, $n \geq 5$, and let $\alpha, a$ be two positive real numbers. Let $f$ be a smooth real-valued function on $M$. We are here concerned with the fourth-order partial differential equation

$$
\begin{equation*}
P_{g} u+a u=f u^{2^{\sharp}-1}, \tag{E}
\end{equation*}
$$

where, as in (0.3),

$$
P_{g} u=\Delta_{g}^{2} u+\alpha \Delta_{g} u
$$

When referring to a solution of $(E)$, we assume that the solution is positive and smooth. Multiplying $(E)$ by $u$, and integrating over $M$, a necessary condition for $(E)$ to have a solution is that $f$ is positive somewhere on $M$. For $u$ in $H_{2}^{2}(M)$, we let

$$
I_{g}(u)=\int_{M}\left(P_{g} u\right) u d v_{g}+a \int_{M} u^{2} d v_{g}
$$

We also let

$$
\mathscr{H}_{f}=\left\{u \in H_{2}^{2}(M) / \int_{M} f|u|^{2^{\sharp}} d v_{g}=1\right\} .
$$

Our first result here is the following theorem.
Theorem 2.1. Let $(M, g)$ be a smooth, compact, $n$-dimensional Riemannian manifold, $n \geq 5$, let $P_{g}$ be the operator given by (0.3), let a be some positive real number, and let $f$ be a smooth positive function on $M$. Then the inequality

$$
\begin{equation*}
\inf _{u \in \mathscr{H}_{f}} I_{g}(u) \leq \frac{1}{\left(\max _{M} f\right)^{2 / 2^{\sharp}} K_{0}} \tag{2.1}
\end{equation*}
$$

always holds, with the additional property that if the inequality in (2.1) is strict, and if $a \leq \alpha^{2} / 4$, then the infimum in the left-hand side of (2.1) is attained by a smooth positive function. In particular, if the inequality in (2.1) is strict and $a \leq \alpha^{2} / 4$, then $(E)$ possesses a smooth positive solution.

Proof. We start by proving (2.1). Suppose on the contrary that

$$
\inf _{u \in \mathscr{H}_{f}} I_{g}(u)>\frac{1}{\left(\max _{M} f\right)^{2 / 2^{\sharp}} K_{0}} .
$$

Then there exists $\epsilon>0$ such that for all $u \in H_{2}^{2}(M)$,

$$
\frac{1}{\left(\max _{M} f\right)^{2 / 2^{\sharp}}}\left(\int_{M} f|u|^{2^{\sharp}} d v_{g}\right)^{2 / 2^{\sharp}} \leq K_{0}(1-\epsilon) \int_{M}\left(P_{g} u\right) u d v_{g}+B\|u\|_{2}^{2},
$$

where, for instance, $B=a K_{0}$. If $x_{0}$ is a point where $f$ is maximum, for $r>0$ sufficiently small, and all $x \in B_{x_{0}}(r)$,

$$
f(x) \geq f\left(x_{0}\right)\left(1-\frac{\epsilon}{2}\right)^{2^{\sharp} / 2} .
$$

Let $\hat{\epsilon}=\epsilon /(2-\epsilon)$ and $\hat{B}=B(1-(\epsilon / 2))^{-1}$. Then for all $u \in \mathscr{D}\left(B_{x_{0}}(r)\right)$,

$$
\|u\|_{2^{\sharp}}^{2} \leq K_{0}(1-\hat{\epsilon}) \int_{M}\left(P_{g} u\right) u d v_{g}+\hat{B}\|u\|_{2}^{2} .
$$

The same arguments as the ones used in the proof of Lemma 1.2 then lead to a contradiction. It follows that (2.1) holds. Let us now prove the second part of the theorem. Let $q \in\left(2,2^{\sharp}\right)$. Set

$$
\mu_{q}=\inf _{u \in \mathscr{H}_{f}^{q}} I_{g}(u)
$$

where

$$
\mathscr{H}_{f}^{q}=\left\{u \in H_{2}^{2}(M) / \int_{M} f|u|^{q} d v_{g}=1\right\} .
$$

Since the embedding of $H_{2}^{2}(M)$ in $L^{q}(M)$ is compact, we know from classical variational arguments that $\mu_{q}$ is attained. In other words, there exists $u_{q} \in \mathscr{H}_{f}^{q}$ such that $I_{g}\left(u_{q}\right)=\mu_{q}$. In particular, $u_{q}$ is a weak solution of

$$
P_{g} u_{q}+a u_{q}=\mu_{q} f\left|u_{q}\right|^{q-2} u_{q}
$$

By classical bootstrap, $u_{q} \in L^{s}(M)$ for all $s$. It easily follows that $u_{q}$ is in fact $C^{3}$. Mimicking what is done in Van der Vorst [23], let $\tilde{u}_{q}$ be the solution of

$$
\Delta_{g} \tilde{u}_{q}+\frac{\alpha}{2} \tilde{u}_{q}=\left|\Delta_{g} u_{q}+\frac{\alpha}{2} u_{q}\right|
$$

Clearly, $\tilde{u}_{q}$ is $C^{2}$, and

$$
\Delta_{g}\left(\tilde{u}_{q} \pm u_{q}\right)+\frac{\alpha}{2}\left(\tilde{u}_{q} \pm u_{q}\right) \geq 0
$$

It follows from the maximum principle that $\tilde{u}_{q} \geq\left|u_{q}\right|$, and that $\tilde{u}_{q}>0$. Noting that

$$
\int_{M}\left(\Delta_{g} \tilde{u}_{q}+\frac{\alpha}{2} \tilde{u}_{q}\right)^{2} d v_{g}=\int_{M}\left(\Delta_{g} u_{q}+\frac{\alpha}{2} u_{q}\right)^{2} d v_{g}
$$

it follows from the assumption $a \leq \alpha^{2} / 4$ that

$$
\begin{aligned}
I_{g}\left(\tilde{u}_{q}\right) & =\int_{M}\left(\Delta_{g} \tilde{u}_{q}+\frac{\alpha}{2} \tilde{u}_{q}\right)^{2} d v_{g}+\left(a-\frac{\alpha^{2}}{4}\right) \int_{M} \tilde{u}_{q}^{2} d v_{g} \\
& =\mu_{q}+\left(a-\frac{\alpha^{2}}{4}\right)\left(\int_{M} \tilde{u}_{q}^{2} d v_{g}-\int_{M} u_{q}^{2} d v_{g}\right) \leq \mu_{q}
\end{aligned}
$$

On the other hand,

$$
\int_{M} f \tilde{u}_{q}^{q} d v_{g} \geq \int_{M} f\left|u_{q}\right|^{q} d v_{g}
$$

Hence,

$$
\hat{u}_{q}=\frac{1}{\left(\int_{M} f \tilde{u}_{q}^{q} d v_{g}\right)^{1 / q}} \tilde{u}_{q}
$$

realises $\mu_{q}$. Here again, $\hat{u}_{q}$ is a solution of

$$
P_{g} \hat{u}_{q}+a \hat{u}_{q}=\mu_{q} f \hat{u}_{q}^{q-1}
$$

By classical regularity, $\hat{u}_{q}$ is in fact $C^{\infty}$. The family $\left(\hat{u}_{q}\right)$ is obviously bounded in $H_{2}^{2}(M)$. Up to the extraction of a subsequence, and for $q \rightarrow 2^{\sharp}$, it converges weakly to some nonnegative $u$ in $H_{2}^{2}(M)$. The embedding $H_{2}^{2}(M) \subset H_{1}^{2}(M)$ being compact, we may also assume that it converges strongly to $u$ in $H_{1}^{2}(M)$. It follows from classical arguments that $u$ is a weak solution of

$$
P_{g} u+a u=\mu f u^{2^{\sharp}-1},
$$

where $\mu$ is given by

$$
\mu=\limsup _{q \rightarrow 2^{\sharp}} \mu_{q} .
$$

By Lemma 2.1 below, $u \in L^{s}(M)$ for all $s$. It easily follows that $u$ is $C^{4}$. From the maximum principle, and noting that

$$
\left(\Delta_{g}+\frac{\alpha}{2}\right)^{2} u \geq 0
$$

we get that $u$ is either positive or the zero function. In both cases, it is actually $C^{\infty}$. Let

$$
\mu_{0}=\inf _{u \in \mathscr{H}_{f}} I_{g}(u) .
$$

It is easily seen that $\mu \leq \mu_{0}$. Coming back to the family $\left(\hat{u}_{q}\right)$, we have that

$$
\begin{aligned}
1 & =\left(\int_{M} f \hat{u}_{q}^{q} d v_{g}\right)^{2 / q} \\
& \leq\left(\max _{M} f\right)^{2 / q}\left(\int_{M} \hat{u}_{q}^{q} d v_{g}\right)^{2 / q} \\
& \leq\left(\max _{M} f\right)^{2 / q} V_{g}^{2\left(1 / q-1 / 2^{\sharp}\right)}\left(\int_{M} \hat{u}_{q}^{2^{\sharp}} d v_{g}\right)^{2 / 2^{\sharp}},
\end{aligned}
$$

where $V_{g}$ stands for the volume of $M$ with respect to $g$. Under the assumption that inequality (2.1) is strict, there exists $\epsilon>0$ such that

$$
\mu_{0}\left(K_{0}+\epsilon\right) \leq \frac{1-\epsilon}{\left(\max _{M} f\right)^{2 / 2^{\sharp}}}
$$

We fix such an $\epsilon$. It follows from Theorem 1.1 that there exists a constant $B$, independent of $q$, such that

$$
\left\|\hat{u}_{q}\right\|_{2^{\sharp}}^{2} \leq\left(K_{0}+\epsilon\right) \int_{M}\left(P_{g} \hat{u}_{q}\right) \hat{u}_{q} d v_{g}+B\left\|\hat{u}_{q}\right\|_{2}^{2} .
$$

Therefore,

$$
1 \leq\left(\max _{M} f\right)^{2 / q} V_{g}^{2\left(1 / q-1 / 2^{\sharp}\right)}\left(K_{0}+\epsilon\right)\left(\mu_{q}+C\left\|\hat{u}_{q}\right\|_{2}^{2}\right),
$$

where $C$ is independent of $q$. As $q \rightarrow 2^{\sharp}$, and since $\mu \leq \mu_{0}$, we get that

$$
\epsilon \leq C\|u\|_{2}^{2}
$$

In particular, $u \not \equiv 0$. It thus follows that $u$ is a smooth positive solution of $(E)$. We are left with the proof that $\mu_{0}$ is attained. As we easily check,

$$
\int_{M} f u^{2^{\sharp}} d v_{g} \leq 1 .
$$

Besides,

$$
\int_{M}\left(P_{g} u\right) u d v_{g}+a \int_{M} u^{2} d v_{g}=\mu \int_{M} f u^{2^{\sharp}} d v_{g}
$$

while according to the definition of $\mu_{0}$,

$$
\int_{M}\left(P_{g} u\right) u d v_{g}+a \int_{M} u^{2} d v_{g} \geq \mu_{0}\left(\int_{M} f u^{2^{\sharp}} d v_{g}\right)^{2 / 2^{\sharp}}
$$

It follows that

$$
\mu\left(\int_{M} f u^{2^{\sharp}} d v_{g}\right)^{1-2 / 2^{\sharp}} \geq \mu_{0} .
$$

Hence, $\mu=\mu_{0}$ and

$$
\int_{M} f u^{2^{\sharp}} d v_{g}=1 .
$$

In particular, $u$ achieves the infimum of the definition of $\mu_{0}$. This ends the proof of the theorem.

The following lemma, based on ideas developed in Van der Vorst [23], has been used in the proof of Theorem 2.1.

Lemma 2.1. Let $(M, g)$ be a smooth, compact, $n$-dimensional Riemannian manifold, let $\alpha$ be a positive real number, let $b$ be a real-valued function defined on $M$, and let $u \in H_{2}^{2}(M)$ be a weak solution of

$$
\Delta_{g}^{2} u+\alpha \Delta_{g} u+\frac{\alpha^{2}}{4} u=b u
$$

If $b \in L^{n / 4}(M)$, then $u \in L^{s}(M)$ for all $s \geq 1$.
Proof. We proceed as in Van der Vorst [23]. As a starting point, we claim that for any $\epsilon>0$, there exists $q_{\epsilon} \in L^{n / 4}(M), f_{\epsilon} \in L^{\infty}(M)$, and a constant $K_{\epsilon}>0$ such that

$$
b u=q_{\epsilon} u+f_{\epsilon}, \quad\left\|q_{\epsilon}\right\|_{n / 4}<\epsilon, \quad\left\|f_{\epsilon}\right\|_{\infty} \leq K_{\epsilon}
$$

Here we may assume that $b \not \equiv 0$, and we let

$$
\begin{aligned}
& \Omega_{k}=\{x \in M /|b|<k\} \\
& \Omega_{l}=\{x \in M /|u|<l\}
\end{aligned}
$$

where if $\hat{\epsilon}$ is such that $(2 \hat{\epsilon})^{4 / n}=\epsilon / 2, k$ and $l$ are chosen such that

$$
\begin{gathered}
\|b\|_{L^{n / 4}\left(M \backslash \Omega_{k}\right)}<\hat{\epsilon} \quad \text { and } \quad\|b\|_{L^{n / 4}\left(M \backslash \Omega_{l}\right)}<\hat{\epsilon}, \\
\Omega_{k} \cap \Omega_{l} \neq \emptyset \quad \text { and } \quad b \not \equiv 0 \quad \text { on } \Omega_{k} \cap \Omega_{l}
\end{gathered}
$$

Given $p \geq 1$ an integer that we fix below, let

$$
q_{\epsilon}= \begin{cases}\frac{1}{p} b & \text { in } \Omega_{k} \cap \Omega_{l} \\ b & \text { in }\left(M \backslash \Omega_{k}\right) \cup\left(M \backslash \Omega_{l}\right),\end{cases}
$$

and

$$
f_{\epsilon}=\left(b-q_{\epsilon}\right) u
$$

Clearly, $f_{\epsilon}=0$ on $M \backslash\left(\Omega_{k} \cap \Omega_{l}\right)$. On the other hand,

$$
\begin{aligned}
\left\|q_{\epsilon}\right\|_{n / 4}^{n / 4} & =\int_{\Omega_{k} \cap \Omega_{l}}\left|q_{\epsilon}\right|^{n / 4} d v_{g}+\int_{M \backslash\left(\Omega_{k} \cap \Omega_{l}\right)}\left|q_{\epsilon}\right|^{n / 4} d v_{g} \\
& \leq \int_{\Omega_{k} \cap \Omega_{l}}\left|q_{\epsilon}\right|^{n / 4} d v_{g}+\int_{M \backslash \Omega_{k}}\left|q_{\epsilon}\right|^{n / 4} d v_{g}+\int_{M \backslash \Omega_{l}}\left|q_{\epsilon}\right|^{n / 4} d v_{g} \\
& \leq\left(\frac{1}{p}\right)^{n / 4} \int_{\Omega_{k} \cap \Omega_{l}}|b|^{n / 4} d v_{g}+2 \hat{\epsilon}
\end{aligned}
$$

so that

$$
\left\|q_{\epsilon}\right\|_{n / 4} \leq \frac{1}{p}\|b\|_{n / 4}+\frac{1}{2} \epsilon
$$

Choosing $p$ such that $p>2\|b\|_{n / 4} / \epsilon$, we get that

$$
\left\|q_{\epsilon}\right\|_{n / 4}<\epsilon
$$

Now, since $f_{\epsilon}=0$ on $M \backslash\left(\Omega_{k} \cap \Omega_{l}\right)$,

$$
\left\|f_{\epsilon}\right\|_{\infty} \leq\left|1-\frac{1}{p}\right| k l
$$

and this proves the above claim. The equation

$$
\Delta_{g}^{2} u+\alpha \Delta_{g} u+\frac{\alpha^{2}}{4} u=b u
$$

may now be written as

$$
L_{g}^{2} u=q_{\epsilon} u+f_{\epsilon}
$$

where

$$
L_{g} u=\Delta_{g} u+\frac{\alpha}{2} u
$$

For any $s>1$ and any $f \in L^{s}(M)$, there exists one and only one $u \in H_{4}^{s}(M)$ such that $L_{g}^{2} u=f$. We let $\mathscr{H}_{\epsilon}$ be the operator

$$
\mathscr{H}_{\epsilon} u=\left(L_{g}\right)^{-2}\left(q_{\epsilon} u\right) .
$$

The preceding equation becomes

$$
u-\mathscr{H}_{\epsilon} u=\left(L_{g}\right)^{-2}\left(f_{\epsilon}\right)
$$

Let $v \in L^{s}(M)$, let $s \geq 2^{\sharp}$, and let $u_{\epsilon}$ be such that

$$
L_{g}^{2} u_{\epsilon}=q_{\epsilon} v
$$

Set $\hat{s}=n s /(n+4 s)$. Clearly, $q_{\epsilon} v \in L^{\hat{s}}(M)$, and it follows from elliptic-type arguments that

$$
\left\|u_{\epsilon}\right\|_{s} \leq C\left\|q_{\epsilon} v\right\|_{\hat{s}}
$$

By Hölder's inequality,

$$
\left\|q_{\epsilon} v\right\|_{\hat{s}} \leq\left\|q_{\epsilon}\right\|_{n / 4}\|v\|_{s}
$$

so that

$$
\left\|u_{\epsilon}\right\|_{s} \leq C \epsilon\|v\|_{s}
$$

In other words, for all $s \geq 2^{\sharp}$, $\mathscr{H}_{\epsilon}$ acts from $L^{s}(M)$ into $L^{s}(M)$, and its norm is less than or equal to $C \epsilon$. Let $s \geq 2^{\sharp}$ be given. For $\epsilon>0$ sufficiently small,

$$
\left\|\mathscr{H}_{\epsilon}\right\|_{L^{s} \rightarrow L^{s}}<\frac{1}{2}
$$

and the operator

$$
\left(I-\mathscr{H}_{\epsilon}\right): L^{s}(M) \longrightarrow L^{s}(M)
$$

has an inverse. Since

$$
\left(I-\mathscr{H}_{\epsilon}\right) u=\left(L_{g}\right)^{-2}\left(f_{\epsilon}\right)
$$

and since $u \in L^{2^{\sharp}}(M)$ and $f_{\epsilon} \in L^{\infty}(M)$, we get that $u \in L^{s}(M)$. The lemma is proved.

On what concerns Theorem 2.1, the equality in (2.1) holds when $(M, g)$ is the standard unit sphere $\left(S^{n}, h\right), P_{g}+a=P_{h}^{n}$, and $f$ is a positive constant. This is just equation (1.9) of Proposition 1.1:

$$
\begin{equation*}
\inf _{u \in C^{\infty}\left(S^{n}\right) \backslash\{0\}} \frac{\int_{S^{n}}\left(P_{h}^{n} u\right) u d v_{h}}{\left(\int_{S^{n}}|u|^{2} d v_{h}\right)^{2 / 2^{\sharp}}}=\frac{1}{K_{0}} . \tag{2.2}
\end{equation*}
$$

Independently, let

$$
u_{0} \equiv\left(\int_{M} f d v_{g}\right)^{-1 / 2^{\sharp}}
$$

Clearly, $u_{0} \in \mathscr{H}_{f}$, and

$$
I_{g}\left(u_{0}\right)=\frac{a V_{g}}{\left(\int_{M} f d v_{g}\right)^{2 / 2^{\sharp}}} .
$$

We then get the following result from Theorem 2.1.
Corollary 2.1. Let $(M, g)$ be a smooth, compact, $n$-dimensional Riemannian manifold, $n \geq 5$, let $P_{g}$ be the operator given by (0.3), let $a>0$ be real, and let $f$ be a smooth, positive function defined on $M$. If $a \leq \alpha^{2} / 4$ and if

$$
\begin{equation*}
\frac{\int_{M} f d v_{g}}{V_{g} \max _{M} f}>\left(a K_{0}\right)^{2^{\sharp} / 2} V_{g}^{2^{\sharp} / 2-1}, \tag{2.3}
\end{equation*}
$$

where $V_{g}$ stands for the volume of $M$ with respect to $g$, then $(E)$ possesses a smooth positive solution.

Here again, the standard unit sphere plays a particular role in this result. As already mentioned in the proof of Proposition 1.1,

$$
\begin{equation*}
d_{n} K_{0}=\omega_{n}^{-4 / n} \tag{2.4}
\end{equation*}
$$

where $\omega_{n}$ is the volume of $S^{n}$ with respect to $h$. It follows that if $(M, g)=\left(S^{n}, h\right)$ and $P_{g}+a=P_{h}^{n}$, then the right-hand side in (2.3) is 1 . On the contrary, the left-hand side is always less than or equal to 1 . The strict inequality (2.3) is therefore never satisfied when $(M, g)=\left(S^{n}, h\right)$ and $P_{g}+a=P_{h}^{n}$. On the other hand, the condition $a \leq \alpha^{2} / 4$ does hold for $P_{h}^{n}$, and we indeed do have that $d_{n} \leq c_{n}^{2} / 4$. As we easily check, the difficulty mentioned above disapears when considering quotients of $S^{n}$. The volume there becomes smaller, and the following result holds. In the particular case $n \leq 7$, see also Theorem 2.2.

Corollary 2.2. Let $\left(S^{n}, h\right)$ be the standard $n$-dimensional unit sphere, $n \geq 5$. For any $\epsilon \in(0,1)$, there exists an integer $k_{\epsilon}$ with the following property: if $f$ smooth on $S^{n}$ is invariant under the action of a subgroup $G$ of $O(n+1)$ acting freely on $S^{n}$ and of order $k \geq k_{\epsilon}$, and if $f$ is such that $\|f-1\|_{C^{0}}<\epsilon$, then the equation

$$
\Delta_{h}^{2} u+c_{n} \Delta_{h} u+d_{n} u=f u^{2^{\sharp}-1}
$$

possesses a smooth, positive, $G$-invariant solution. In particular, there exists a metric $g$ in the conformal class of $h$ for which $Q_{g}^{n}=f$, where

$$
Q_{g}^{n}=\frac{1}{2(n-1)} \Delta_{g} S_{g}+\frac{n^{3}-4 n^{2}+16 n-16}{8(n-1)^{2}(n-2)^{2}} S_{g}^{2}-\frac{2}{(n-2)^{2}}\left|R c_{g}\right|^{2},
$$

and for which $G \subset \operatorname{Isom}_{g}\left(S^{n}\right)$, where $\operatorname{Isom}_{g}\left(S^{n}\right)$ stands for the isometry group of $S^{n}$ with respect to $g$.

Proof. Let $M$ be the quotient manifold $S^{n} / G$, and let $g_{0}$ be its standard metric induced by $h$. We still denote by $f$ the quotient of $f$ on $M$. The Paneitz-Branson operator on $M$ is given by

$$
P_{g_{0}}^{n} u=\Delta_{g_{0}}^{2} u+c_{n} \Delta_{g_{0}} u+d_{n} u,
$$

and as already mentioned, $d_{n} \leq c_{n}^{2} / 4$. If $f$ is such that $\|f-1\|_{C^{0}}<\epsilon$, then

$$
\frac{\int_{M} f d v_{g_{0}}}{V_{g_{0}} \max _{M} f}>\frac{1-\epsilon}{1+\epsilon}
$$

Now,

$$
\left(d_{n} K_{0}\right)^{2^{\sharp} / 2} V_{g_{0}}^{2^{\sharp} / 2-1}=\frac{1}{k^{2^{\sharp} / 2-1}}\left(\left(d_{n} K_{0}\right)^{2^{\sharp} / 2} \omega_{n}^{2^{\sharp} / 2-1}\right) .
$$

Then set

$$
k_{\epsilon}=\left[\left(\frac{1+\epsilon}{1-\epsilon}\right)^{2 /\left(2^{\sharp}-2\right)}\left(d_{n} K_{0}\right)^{2^{\sharp} /\left(2^{\sharp}-2\right)} \omega_{n}\right]+1,
$$

where $[x]$ stands for the largest integer not exceeding $x$. If $k \geq k_{\epsilon}$, then

$$
\frac{\int_{M} f d v_{g_{0}}}{V_{g_{0}} \max _{M} f}>\left(d_{n} K_{0}\right)^{2^{\sharp} / 2} V_{g_{0}}^{2^{\sharp} / 2-1} .
$$

Noticing that the existence of a solution to the equation on $M$ gives the existence of a $G$-invariant solution of the equation on $S^{n}$, the result follows from Theorem 2.1.

We now concentrate on the study of $(E)$ when $(M, g)$ is the standard sphere and $P_{g}+a=P_{h}^{n}$. The following result, together with Theorem 2.3, shows that $c_{n}$ in the definition of $P_{h}^{n}$ is critical. In the study of $(E)$ on the standard sphere, we do get obstructions by Theorem 2.3 when $\alpha=c_{n}$ and $a=d_{n}$. These obstructions disappear according to Corollary 2.3 if $\alpha<c_{n}$. When studying $(E)$ on the standard sphere, both the medium term $\alpha=c_{n}$ and the nonlinear growth $p=2^{\sharp}-1$ are critical. For more details on such assertions, we refer to the remark after the proof of Theorem 2.3.

Corollary 2.3. Let $\left(S^{n}, h\right)$ be the standard $n$-dimensional unit sphere, $n \geq 6$, let $\alpha$ and a be two positive real numbers, and let $f$ be a smooth, positive function on $S^{n}$. If $a \leq \alpha^{2} / 4, \alpha<c_{n}$, and if

$$
(n-6) \frac{\Delta_{h} f\left(x_{0}\right)}{f\left(x_{0}\right)}<\frac{8 n(n-1)}{(n-4)(n+2)}\left(c_{n}-\alpha\right)
$$

for at least one $x_{0}$ where $f$ is maximum, then the equation

$$
\Delta_{h}^{2} u+\alpha \Delta_{h} u+a u=f u^{2^{\sharp}-1}
$$

possesses a smooth, positive solution.
Proof. Let $x_{0}$ be a point where $f$ is maximum, and $r$ be the distance on $S^{n}$ to $x_{0}$. For $\beta>1$, we let $u_{\beta}$ be the function

$$
u_{\beta}=\frac{\left(\beta^{2}-1\right)^{(n-4) / 4}}{(\beta-\cos r)^{(n-4) / 2}}
$$

As already mentioned in the proof of Proposition 1.1, $h$ and $u_{\beta}^{4 /(n-4)} h$ are isometric. It follows that

$$
P_{h}^{n} u_{\beta}=d_{n} u_{\beta}^{2^{\sharp}-1}
$$

where $P_{h}^{n}$ is the Paneitz-Branson operator on the sphere, as defined in (0.1). According to the developments made in the proof of Proposition 1.2,

$$
\begin{aligned}
\int_{S^{n}} & \left(P_{h} u_{\beta}\right) u_{\beta} d v_{h}+a \int_{S^{n}} u_{\beta}^{2} d v_{h} \\
& =d_{6} \omega_{6}+A\left(\alpha-c_{6}\right)(\beta-1) \ln \left(\frac{1}{\beta-1}\right)+o\left((\beta-1) \ln \left(\frac{1}{\beta-1}\right)\right) \quad \text { if } n=6 \\
& =d_{n} \omega_{n}+2^{n-3}(n-4)^{2}\left(\alpha-c_{n}\right)(\beta-1) \omega_{n-1} I+o(\beta-1) \quad \text { if } n>6,
\end{aligned}
$$

where $A$ is some positive constant and

$$
I=\int_{0}^{+\infty} \frac{y^{n+1} d y}{\left(1+y^{2}\right)^{n-2}}
$$

We now write

$$
f=f\left(x_{0}\right)+(1-\cos r) \hat{f}
$$

It is easily seen that

$$
\lim _{t \rightarrow 0^{+}} \frac{\int_{\partial B_{x_{0}}(t)} \hat{f} d \sigma}{V_{h}\left(\partial B_{x_{0}}(t)\right)}=-\frac{\omega_{n-1}}{n} \Delta_{h} f\left(x_{0}\right)
$$

where $V_{h}\left(\partial B_{x_{0}}(t)\right)$ stands for the area of $\partial B_{x_{0}}(t)$ with respect to the metric induced by $h$. By the changes of variables $x=\tan (r / 2)$, and then $y=\sqrt{(\beta+1) /(\beta-1)} x$, we get that

$$
\begin{aligned}
(\beta-1)^{-1} & \int_{B_{x_{0}}(t)}(1-\cos r) u_{\beta}^{2^{\sharp}} d v_{h} \\
& =\frac{2^{n+1} \omega_{n-1}}{\beta+1} \int_{0}^{\sqrt{(\beta+1) /(\beta-1)} T} \frac{y^{n+1} d y}{\left(1+((\beta-1) /(\beta+1)) y^{2}\right)\left(1+y^{2}\right)^{n}}
\end{aligned}
$$

for all $t \in(0, \pi)$, where $T=\tan (t / 2)$. It easily follows that

$$
\lim _{\beta \rightarrow 1^{+}}(\beta-1) \int_{S^{n}}(1-\cos r) \hat{f} u_{\beta}^{2^{\sharp}} d v_{h}=-\frac{2^{n} \omega_{n-1}}{n} J \Delta_{h} f\left(x_{0}\right),
$$

where

$$
J=\int_{0}^{+\infty} \frac{y^{n+1} d y}{\left(1+y^{2}\right)^{n}}
$$

As a consequence,

$$
\int_{S^{n}} f u_{\beta}^{2^{\sharp}} d v_{h}=f\left(x_{0}\right) \omega_{n}\left(1-\frac{2^{n} \omega_{n-1}}{n \omega_{n}}(\beta-1) J \frac{\Delta_{h} f\left(x_{0}\right)}{f\left(x_{0}\right)}+o(\beta-1)\right) .
$$

Since $d_{n} K_{0}=\omega_{n}^{-4 / n}$, and since $x_{0}$ is a point where $f$ is maximum, we get that

$$
\begin{aligned}
& \frac{\int_{S^{n}}\left(P_{h} u_{\beta}\right) u_{\beta} d v_{h}+a \int_{S^{n}} u_{\beta}^{2} d v_{h}}{} \\
& \quad\left(\int_{S^{n}} f u_{\beta}^{2^{\sharp}} d v_{h}\right)^{2 / 2^{\sharp}} \\
& \quad=\frac{1}{K_{0}\left(\max _{M} f\right)^{2 / 2^{\sharp}}}\left(1+B\left(\alpha-c_{6}\right) \varepsilon_{\beta}+o\left(\varepsilon_{\beta}\right)\right) \quad \text { if } n=6 \\
& \quad=\frac{1}{K_{0}\left(\max _{M} f\right)^{2 / 2^{\sharp}}}(1+C(\beta-1)+o(\beta-1)) \quad \text { if } n>6,
\end{aligned}
$$

where $\varepsilon_{\beta}=(\beta-1) \ln (1 /(\beta-1)), B>0$ does not depend on $\beta$, and

$$
C=\frac{2^{n+1} \omega_{n-1}}{2^{\sharp} n \omega_{n}} J\left(\frac{\Delta_{h} f\left(x_{0}\right)}{f\left(x_{0}\right)}+\frac{2 n I}{\left(n^{2}-4\right) J}\left(\alpha-c_{n}\right)\right) .
$$

As we easily check, see, for instance, Demengel-Hebey [8],

$$
I=\frac{1}{2} \frac{\Gamma((n+2) / 2) \Gamma((n-6) / 2)}{\Gamma(n-2)} \quad \text { and } \quad J=\frac{1}{2} \frac{\Gamma((n+2) / 2) \Gamma((n-2) / 2)}{\Gamma(n)}
$$

Hence,

$$
\frac{2 n I}{\left(n^{2}-4\right) J}=\frac{8 n(n-1)}{(n-6)(n-4)(n+2)}
$$

Under our assumptions $C<0$, so that for $\beta>1$ sufficiently close to 1 ,

$$
\frac{\int_{S^{n}}\left(P_{h} u_{\beta}\right) u_{\beta} d v_{h}+a \int_{S^{n}} u_{\beta}^{2} d v_{h}}{\left(\int_{S^{n}} f u_{\beta}^{2 \sharp} d v_{h}\right)^{2 / 2^{\sharp}}}<\frac{1}{K_{0}\left(\max _{M} f\right)^{2 / 2^{\sharp}}}
$$

The result now follows from Theorem 2.1.
The equation involved in the study of the prescribed scalar curvature problem on the sphere, also referred to as the Kazdan-Warner problem or the Nirenberg problem, is the equation

$$
\Delta_{h} u+\frac{n(n-2)}{4} u=f u^{2^{\star}-1}
$$

where $2^{\star}=2 n /(n-2)$. In the study of this equation, a celebrated result of Escobar and Schoen [11] states that if $n=3$ and if $f$ is invariant under the action of a nontrivial subgroup $G$ of $O(4)$ acting freely on $S^{3}$, then the above equation possesses a smooth, positive $G$-invariant solution. In particular, under these assumptions, $f$ is the scalar curvature of a $G$-invariant conformal metric to $h$. The same result was proved by Moser [20] when $n=2$ and $f$ is assumed to be invariant under the action of the antipodal group $G=\{I d,-I d\}$, the only group acting freely on $S^{n}$ when the dimension $n$ is even. A natural question is whether such types of results do hold for the equation

$$
P_{h}^{n} u=f u^{2^{\sharp}-1}
$$

This is the subject of the following theorem. As a first remark, note that by Edmunds-Fortunato-Janelli [10] and Pucci-Serrin [22], low dimensions for the Euclidean biharmonic operator are $n=5,6,7$. As another remark, we mention that there should be an analogue of our result when $G$ acts without fixed points (i.e., for any $x$, the $G$-orbit of $x$ has at least two elements). Concerning the above mentioned scalar curvature problem on the sphere, this was proved by Hebey [13].

Theorem 2.2. Let $\left(S^{n}, h\right)$ be the standard $n$-dimensional unit sphere, $n=5,6$, or 7, and let $f$ be a smooth positive function on $S^{n}$. We assume that $f$ is invariant under the action of a nontrivial subgroup $G$ of $O(n+1)$ acting freely on $S^{n}$, and if $n=6$ or 7 , we assume that $\Delta_{h} f(x)=0$ for at least one $x$ where $f$ is maximum. Then the equation

$$
P_{h}^{n} u=f u^{2^{\sharp}-1}
$$

possesses a smooth, positive G-invariant solution. In particular, there exists a metric $g$ in the conformal class of $h$ for which $Q_{g}^{n}=f$, where

$$
Q_{g}^{n}=\frac{1}{2(n-1)} \Delta_{g} S_{g}+\frac{n^{3}-4 n^{2}+16 n-16}{8(n-1)^{2}(n-2)^{2}} S_{g}^{2}-\frac{2}{(n-2)^{2}}\left|R c_{g}\right|^{2}
$$

and for which $G \subset \operatorname{Isom}_{g}\left(S^{n}\right)$, where $\operatorname{Isom}_{g}\left(S^{n}\right)$ stands for the isometry group of $S^{n}$ with respect to $g$.

Proof. Let $M$ be the quotient manifold $S^{n} / G$, and let $g_{0}$ be its standard metric induced by $h$. Also let $u \in H_{2}^{2}(M), u \not \equiv 0$, and let $\tilde{u}$ be the function on $S^{n}$ induced by $u$. As we easily check,

$$
\frac{\int_{S^{n}}\left(P_{h}^{n} \tilde{u}\right) \tilde{u} d v_{h}}{\left(\int_{S^{n}} f|\tilde{u}|^{2^{\sharp}} d v_{h}\right)^{2 / 2^{\sharp}}}=k^{1-2 / 2^{\sharp}} \frac{\int_{M}\left(P_{g_{0}}^{n} \tilde{u}\right) \tilde{u} d v_{g_{0}}}{\left(\int_{M} f|\tilde{u}|^{2^{\sharp}} d v_{g_{0}}\right)^{2 / 2^{\sharp}}},
$$

where $k$ is the number of elements in $G$, and $f$ in the right-hand side of this relation stands for the quotient of $f$ on $M$. The existence of a solution to the equation on $M$ leads to the existence of a $G$-invariant solution to the equation on $S^{n}$. Therefore, as a consequence of Theorem 2.1, it suffices to show that

$$
\begin{equation*}
\inf _{u \in \Lambda} \frac{\int_{S^{n}}\left(P_{h}^{n} u\right) u d v_{h}}{\left(\int_{S^{n}} f|u|^{\sharp} d v_{h}\right)^{2 / 2^{\sharp}}}<\frac{k^{1-2 / 2^{\sharp}}}{\left(\max _{S^{n}} f\right)^{2 / 2^{\sharp}} K_{0}}, \tag{2.5}
\end{equation*}
$$

where $\Lambda$ stands for the subset of $H_{2}^{2}\left(S^{n}\right)$ consisting of nonzero $G$-invariant functions. Now let $x_{1}$ be a point where $f$ is maximum, and denote by

$$
O_{G}\left(x_{1}\right)=\left\{x_{1}, \ldots, x_{k}\right\}
$$

the $G$-orbit of $x_{1}$. If $r_{i}$ stands for the distance on $S^{n}$ to $x_{i}$, let $u_{i, \beta}, \beta>1$ be the functions on $S^{n}$ defined by

$$
u_{i, \beta}=\frac{\left(\beta^{2}-1\right)^{(n-4) / 4}}{\left(\beta-\cos r_{i}\right)^{(n-4) / 2}}
$$

As already mentioned, $h$ and $u_{i, \beta}^{4 /(n-4)} h$ are isometric. In particular,

$$
P_{h}^{n} u_{i, \beta}=d_{n} u_{i, \beta}^{2 \sharp-1}
$$

and

$$
\int_{S^{n}} u_{i, \beta}^{2^{\sharp}} d v_{h}=\omega_{n} .
$$

Then let

$$
u_{\beta}=\sum_{i=1}^{k} u_{i, \beta}
$$

On the one hand, $u_{\beta}$ is $G$-invariant. On the other hand,

$$
\int_{S^{n}}\left(P_{h}^{n} u_{\beta}\right) u_{\beta} d v_{h}=k d_{n} \omega_{n}+k d_{n} \int_{S^{n}} u_{1, \beta}^{2^{\sharp}-1}\left(\sum_{i=2}^{k} u_{i, \beta}\right) d v_{h} .
$$

Set

$$
A(n)=2^{1+3 n / 4} \omega_{n-1} \int_{0}^{+\infty} \frac{y^{n-1} d y}{\left(1+y^{2}\right)^{(n+4) / 2}}
$$

We claim that for all $t$ in $(0, \pi)$,

$$
\lim _{\beta \rightarrow 1^{+}}(\beta-1)^{1-n / 4} \int_{B_{x_{1}}(t)} u_{1, \beta}^{2^{\sharp}-1} d v_{h}=A(n)
$$

Indeed, by the change of variables $x=\tan \left(r_{1} / 2\right)$ and $y=\sqrt{(\beta+1) /(\beta-1)} x$, we get that

$$
\begin{aligned}
\int_{B_{x_{1}}(t)} & u_{1, \beta}^{2^{\sharp}-1} d v_{h} \\
= & C_{1}(n, \beta)(\beta-1)^{n / 4-1} \\
& \times \int_{0}^{\sqrt{(\beta+1) /(\beta-1)} T} \frac{y^{n-1} d y}{\left(1+((\beta-1) /(\beta+1)) y^{2}\right)^{(n-4) / 2}\left(1+y^{2}\right)^{(n+4) / 2}}
\end{aligned}
$$

where

$$
C_{1}(n, \beta)=2^{n}(\beta+1)^{1-n / 4} \omega_{n-1}
$$

and $T=\tan (t / 2)$. The above claim then easily follows. With easier arguments, for all $h \in C^{0}\left(S^{n}\right)$, and all $t$ in $(0, \pi)$,

$$
\lim _{\beta \rightarrow 1^{+}}(\beta-1)^{2-n / 2} \int_{S^{n} \backslash B_{x_{1}}(t)} h u_{1, \beta}^{2^{\sharp}-1}\left(\sum_{i=2}^{k} u_{i, \beta}\right) d v_{h}=0 .
$$

It follows from these two relations that for all $h \in C^{0}\left(S^{n}\right)$, all $t$ in $(0, \pi)$, and all open subset $\Omega$ of $S^{n}$ that contains $x_{1}$,

$$
\begin{equation*}
\lim _{\beta \rightarrow 1^{+}}(\beta-1)^{2-n / 2} \int_{\Omega} h u_{1, \beta}^{2^{\sharp}-1}\left(\sum_{i=2}^{k} u_{i, \beta}\right) d v_{h}=A(n)\left(\sum_{i=2}^{k} \tilde{u}_{i, 1}\left(x_{1}\right)\right) h\left(x_{1}\right), \tag{2.6}
\end{equation*}
$$

where $\tilde{u}_{i, 1}=\left(1-\cos r_{i}\right)^{2-n / 2}$. Then let $t_{0}>0$ be such that $B_{x_{i}}\left(t_{0}\right) \cap B_{x_{j}}\left(t_{0}\right)=\emptyset$ for $i \neq j$. Since

$$
(a+b)^{2^{\sharp}} \geq a^{2^{\sharp}}+2^{\sharp} a^{2^{\sharp}-1} b,
$$

we may write that

$$
\begin{aligned}
\int_{S^{n}} f u_{\beta}^{2^{\sharp}} d v_{h} & \geq \sum_{i=1}^{k} \int_{B_{x_{i}}\left(t_{0}\right)} f u_{\beta}^{2^{\sharp}} d v_{h} \\
& \geq \sum_{i=1}^{k} \int_{B_{x_{i}\left(t_{0}\right)}} f u_{i, \beta}^{2^{\sharp}} d v_{h}+2^{\sharp} \sum_{i=1}^{k} \sum_{j \neq i} \int_{B_{x_{i}}\left(t_{0}\right)} f u_{i, \beta}^{2^{\sharp}-1} u_{j, \beta} d v_{h} \\
& =k \int_{B_{x_{1}\left(t_{0}\right)}} f u_{1, \beta}^{2^{\sharp}} d v_{h}+2^{\sharp} k \int_{B_{x_{1}}\left(t_{0}\right)} f u_{1, \beta}^{2^{\sharp}-1}\left(\sum_{i=2}^{k} u_{i, \beta}\right) d v_{h} .
\end{aligned}
$$

It follows that

$$
\begin{aligned}
& \int_{S^{n}} f u_{\beta}^{2^{\sharp}} d v_{h} \geq k f\left(x_{1}\right) \omega_{n}-k f\left(x_{1}\right) \int_{S^{n} \backslash B_{x_{1}}\left(t_{0}\right)} u_{1, \beta}^{2^{\sharp}} d v_{h} \\
& \quad+k \int_{B_{x_{1}}\left(t_{0}\right)}\left(f-f\left(x_{1}\right)\right) u_{1, \beta}^{2^{\sharp}} d v_{h}+2^{\sharp} k \int_{B_{x_{1}}\left(t_{0}\right)} f u_{1, \beta}^{2^{\sharp}-1}\left(\sum_{i=2}^{k} u_{i, \beta}\right) d v_{h} .
\end{aligned}
$$

It is easily seen that

$$
\begin{equation*}
\lim _{\beta \rightarrow 1^{+}}(\beta-1)^{2-n / 2} \int_{S^{n} \backslash B_{x_{1}}\left(t_{0}\right)} u_{1, \beta}^{2^{\sharp}} d v_{h}=0 . \tag{2.7}
\end{equation*}
$$

On the other hand, for any $5 \leq n \leq 7$,

$$
\begin{equation*}
\lim _{\beta \rightarrow 1^{+}}(\beta-1)^{2-n / 2} \int_{B_{x_{1}}\left(t_{0}\right)}\left(f-f\left(x_{1}\right)\right) u_{1, \beta}^{2^{\sharp}} d v_{h}=0 . \tag{2.8}
\end{equation*}
$$

Indeed, suppose that $n=5$. Since $x_{1}$ is a critical point for $f$, there exists a constant $C>0$ such that for all $x \in B_{x_{1}}\left(t_{0}\right)$,

$$
\left|f(x)-f\left(x_{1}\right)\right| \leq C\left(1-\cos r_{1}\right)
$$

With the change of variables $x=\tan \left(r_{1} / 2\right)$ and $y=\sqrt{(\beta+1) /(\beta-1)} x$, we get that

$$
(\beta-1)^{2-n / 2} \int_{B_{x_{1}}\left(t_{0}\right)}\left(1-\cos r_{1}\right) u_{1, \beta}^{2^{\sharp}} d v_{h}=O\left((\beta-1)^{3-n / 2}\right),
$$

from which (2.8) follows. Suppose then that $n=6$ or 7 and that $x_{1}$ is such that $\Delta_{h} f\left(x_{1}\right)=0$. We may write that there exists a constant $C>0$ such that for all $x \in B_{x_{1}}\left(t_{0}\right)$,

$$
\left|f(x)-f\left(x_{1}\right)\right| \leq C\left(1-\cos r_{1}\right)^{2}
$$

As above, we get that

$$
(\beta-1)^{2-n / 2} \int_{B_{x_{1}}\left(t_{0}\right)}\left(1-\cos r_{1}\right)^{2} u_{1, \beta}^{2^{\sharp}} d v_{h}=O\left((\beta-1)^{4-n / 2}\right),
$$

from which (2.8) also follows. Now, by (2.4), (2.6), (2.7), and (2.8),

$$
\frac{\int_{S^{n}}\left(P_{h}^{n} u_{\beta}\right) u_{\beta} d v_{h}}{\left(\int_{S^{n}} f u_{\beta}^{2^{\sharp}} d v_{h}\right)^{2 / 2^{\sharp}}}=\frac{k^{1-2 / 2^{\sharp}}}{f\left(x_{1}\right)^{2 / 2^{\sharp}} K_{0}} \times \frac{1+(\beta-1)^{n / 2-2} A_{k}(n)+o\left((\beta-1)^{n / 2-2}\right)}{1+2(\beta-1)^{n / 2-2} A_{k}(n)+o\left((\beta-1)^{n / 2-2}\right)},
$$

where $A_{k}(n)>0$ is given by

$$
A_{k}(n)=\frac{A(n)}{\omega_{n}} \sum_{i=2}^{k} \tilde{u}_{i, 1}\left(x_{1}\right)
$$

Hence, for every $\beta>1$ sufficiently close to 1 ,

$$
\frac{\int_{S^{n}}\left(P_{h}^{n} u_{\beta}\right) u_{\beta} d v_{h}}{\left(\int_{S^{n}} f u_{\beta}^{2^{\sharp}} d v_{h}\right)^{2 / 2^{\sharp}}}<\frac{k^{1-2 / 2^{\sharp}}}{f\left(x_{1}\right)^{2 / 2^{\sharp}} K_{0}} .
$$

In particular, since $u_{\beta}$ is $G$-invariant, and $f\left(x_{1}\right)=\max _{S^{n}} f$,

$$
\inf _{u \in \Lambda} \frac{\int_{S^{n}}\left(P_{h}^{n} u\right) u d v_{h}}{\left(\int_{S^{n}} f|u|^{2^{\sharp}} d v_{h}\right)^{2 / 2^{\sharp}}}<\frac{k^{1-2 / 2^{\sharp}}}{\left(\max _{S^{n}} f\right)^{2 / 2^{\sharp}} K_{0}} .
$$

This is exactly inequality (2.5). The theorem is thus proved.
A celebrated result of Kazdan and Warner [17] states that the scalar curvature equation on the sphere $\left(S^{n}, h\right)$ possesses obstructions. We prove here that such obstructions hold similarly for the equation

$$
P_{h}^{n} u=f u^{2^{\sharp}-1} .
$$

In the statement of Theorem 2.3, $(\nabla f \nabla \varphi)$ stands for the pointwise scalar product with respect to $h$ of $\nabla f$ and $\nabla \varphi$.

Theorem 2.3. Let $\left(S^{n}, h\right)$ be the standard $n$-dimensional unit sphere, $n \geq 5$, and let $f$ be a smooth function on $S^{n}$, positive somewhere on $S^{n}$. If $u$ is a smooth positive solution of the equation

$$
\begin{equation*}
P_{h}^{n} u=f u^{2^{\sharp}-1}, \tag{2.9}
\end{equation*}
$$

where $P_{h}^{n}$ is as in (0.1), then for any eigenfunction $\varphi$ of $\Delta_{h}$ associated to the first nonzero eigenvalue $\lambda_{1}=n$,

$$
\int_{S^{n}}(\nabla f \nabla \varphi) u^{2^{\sharp}} d v_{h}=0 .
$$

In particular, for any $\epsilon>0$ and any eigenfunction $\varphi \not \equiv 0$ of $\Delta_{h}$ associated to the first nonzero eigenvalue $\lambda_{1}=n$, (2.9) with $f=1+\epsilon \varphi$ does not possess a smooth, positive solution.

Proof. The proof mainly follows what was done in Kazdan and Warner [17]. Let $\varphi$ be an eigenfunction of $\Delta_{h}$ associated to the first nonzero eigenvalue $\lambda_{1}=n$, and let $u$ be a smooth function on $S^{n}$. As it is easy to see,

$$
\left(\Delta_{h}^{2} u\right)(\nabla u \nabla \varphi) \simeq\left(\Delta_{h} u\right)\left(\nabla \Delta_{h} u \nabla \varphi\right)-(n-2)\left(\Delta_{h} u\right)(\nabla u \nabla \varphi)-2\left(\Delta_{h} u\right)^{2} \varphi
$$

where the sign " $\simeq$ " means that the relation holds modulo terms in divergence form. Clearly,

$$
\left(\Delta_{h} u\right)\left(\nabla \Delta_{h} u \nabla \varphi\right) \simeq \frac{1}{2} n \varphi\left(\Delta_{h} u\right)^{2}
$$

so that

$$
\left(\Delta_{h} u\right)^{2}(\nabla u \nabla \varphi) \simeq \frac{n-4}{2}\left(\Delta_{h} u\right)^{2} \varphi-(n-2)\left(\Delta_{h} u\right)(\nabla u \nabla \varphi)
$$

Suppose now that $u$ is a solution of (2.9). Then

$$
\left(\Delta_{h} u\right)^{2} \varphi \simeq 2\left(\Delta_{h} u\right)(\nabla u \nabla \varphi)-n u \varphi\left(\Delta_{h} u\right)+f \varphi u^{2^{\sharp}}-d_{n} u^{2} \varphi-c_{n} u \varphi\left(\Delta_{h} u\right)
$$

so that

$$
\begin{aligned}
&\left(c_{n}-2\right)\left(\Delta_{h} u\right)(\nabla u \nabla \varphi)-\frac{n-4}{2}\left(n+c_{n}\right) u \varphi\left(\Delta_{h} u\right)+d_{n} u(\nabla u \nabla \varphi) \\
&+\frac{n-4}{2} f \varphi u^{2^{\sharp}}-\frac{n-4}{2} d_{n} u^{2} \varphi \simeq f u^{2^{\sharp}-1}(\nabla u \nabla \varphi) .
\end{aligned}
$$

Since

$$
\left(\Delta_{h} u\right)(\nabla u \nabla \varphi) \simeq-\frac{n-2}{2}\left(\frac{1}{2} n u^{2} \varphi-u \varphi\left(\Delta_{h} u\right)\right)
$$

and

$$
\frac{n-2}{2}\left(c_{n}-2\right)=\frac{n-4}{2}\left(n+c_{n}\right)
$$

the terms $u \varphi\left(\Delta_{h} u\right)$ disappear. We then get that
$-\frac{n-2}{4}\left(c_{n}-2\right) n u^{2} \varphi+d_{n} u(\nabla u \nabla \varphi)+\frac{n-4}{2} f \varphi u^{2^{\sharp}}-\frac{n-4}{2} d_{n} u^{2} \varphi \simeq f u^{2^{\sharp}-1}(\nabla u \nabla \varphi)$.

Now it is easily seen that

$$
u(\nabla u \nabla \varphi) \simeq \frac{n}{2} u^{2} \varphi
$$

and that

$$
f u^{2^{\sharp}-1}(\nabla u \nabla \varphi) \simeq-\frac{1}{2^{\sharp}} u^{2^{\sharp}}(\nabla f \nabla \varphi)+\frac{n}{2^{\sharp}} f u^{2^{\sharp}} \varphi .
$$

Therefore,

$$
\left(2 d_{n}-\frac{n(n-2)}{4}\left(c_{n}-2\right)\right) u^{2} \varphi+\frac{1}{2^{\sharp}} u^{2^{\sharp}}(\nabla f \nabla \varphi) \simeq 0 .
$$

Since

$$
2 d_{n}=\frac{n(n-2)}{4}\left(c_{n}-2\right)
$$

we find that

$$
u^{2^{\sharp}}(\nabla f \nabla \varphi) \simeq 0 .
$$

This ends the proof of the theorem.
To conclude, we collect a few remarks on Corollary 2.3 and Theorem 2.3. First, let $\varphi$ be an eigenfunction of $\Delta_{h}$ associated to the first nonzero eigenvalue $\lambda_{1}=n$, and, for $\varepsilon>0$, set $f_{\varepsilon}=1+\varepsilon \varphi$. Given $\alpha>0$, consider the equation

$$
\Delta_{h}^{2} u+\alpha \Delta_{h} u+d_{n} u=f_{\varepsilon} u^{2^{\sharp}-1} .
$$

According to Corollary 2.3, if $n>6$ and $\alpha \in\left[2 \sqrt{d_{n}}, c_{n}\right)$, then there exists $\varepsilon_{\alpha}>0$ such that if $\varepsilon \leq \varepsilon_{\alpha},\left(E_{\alpha}^{\varepsilon}\right)$ possesses a smooth positive solution. On the contrary, by Theorem 2.3, for all $n \geq 5$ and all $\varepsilon>0,\left(E_{c_{n}}^{\varepsilon}\right)$ does not possess any smooth positive solution. This is one of the possible illustrations of the criticality of $c_{n}$ we mentioned before stating Corollary 2.3. As another remark, note that Theorem 2.3, together with Theorem 2.1, gives another proof of (1.9). Indeed, suppose by contradiction that (1.9) is false. In other words, assume that

$$
\inf _{u \in C^{\infty}\left(S^{n}\right) \backslash\{0\}} \frac{\int_{S^{n}}\left(P_{h}^{n} u\right) u d v_{h}}{\left(\int_{S^{n}}|u|^{2^{\sharp}} d v_{h}\right)^{2 / 2^{\sharp}}}<\frac{1}{K_{0}} .
$$

Then for any $f$ sufficiently close to 1 in the $C^{0}$-topology,

$$
\inf _{u \in C^{\infty}\left(S^{n}\right) \backslash\{0\}} \frac{\int_{S^{n}}\left(P_{h}^{n} u\right) u d v_{h}}{\left(\int_{S^{n}} f|u|^{2^{\sharp}} d v_{h}\right)^{2 / 2^{\sharp}}}<\frac{1}{\left(\max _{S^{n}} f\right)^{2 / 2^{\sharp}} K_{0}} .
$$

It follows from Theorem 2.1 that for such an $f$, the equation

$$
P_{h}^{n} u=f u^{2^{\sharp}-1}
$$

has a smooth, positive solution. This is in contradiction to the last part of Theorem 2.3 and thus proves (1.9).

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