

# Chap. 4: Flocking of kinetic equations

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In this last chapter, we want to prove flocking on the kinetic equation - as we already performed at the microscopic level.

We now choose  $\Psi(s)$  such that  $\exists \Psi_m > 0$  such that  $\Psi(s) \geq \Psi_m \quad \forall s \in \mathbb{R}^+$

We lose the integrability condition and have long range interactions.

We consider the (CS) model

$$\begin{cases} \partial_t f + v \nabla_x f + \operatorname{div} (L[f] f) + \beta \operatorname{div} ((u-v) f) = 0 \\ L[f] = \Psi * \rho u - (\Psi * \rho) v \\ \Psi \text{ is symmetric} \end{cases}$$

We define the following Lyapunov functionals

$$E_1(t) = \int_{\mathbb{R}^d \times \mathbb{R}^d} f |v-u|^2 dx dv \quad \text{relative energy}$$

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and

$$E_2(t) = \int_{\mathbb{R}^{2d}} \rho(t,x) \rho(t,y) |u(t,x) - u(t,y)|^2 dx dy$$

$E_1$  measures the local alignment whereas  
 $E_2$  measures alignment of the macroscopic velocities

We set 
$$E(t) = E_1(t) + \frac{1}{2} E_2(t)$$

and prove the following theorem

Theorem 
$$E(t) \leq E(0) e^{-ct}$$

where 
$$c = 2 \min(\tau, \Psi_m)$$

and 
$$\Psi_m = \min_{x,y} \Psi(x-y) > 0$$

Let us give the main steps for the proof of this result

$$\frac{d\varepsilon_1}{dt} = \underbrace{2 \int f(u-v) \partial_t u + \int \partial_t f |u-v|^2 dx dv}_{=0} \quad (3)$$

$$= 2 \int \nabla_x u \cdot (u-v) \cdot v f \, dx dv = I_1$$

$$- 2 \int [(u-v) \cdot L[f]] f \, dx dv = \underline{I_2} \leq 0$$

$$- \underbrace{2 \int |u-v|^2 f \, dx dv}_{2\varepsilon_1}$$

$$I_1 = 2 \int \operatorname{div} P \cdot u \, dx \quad \text{where}$$

$$P \text{ is the stress tensor } P = \int (v-u) \otimes (v-u) f \, dv$$

Hence we get that

$$\frac{d\varepsilon_1}{dt} \leq 2 \int \operatorname{div}_x P \cdot u \, dx - 2\varepsilon_1$$


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We next estimate

$$\frac{d\varepsilon_2}{dt} = 2 \int \partial_t \rho(t,x) \rho(t,y) \overset{J_1}{|u(t,x) - u(t,y)|^2} dx dy$$

$$+ 2 \int \rho(t,x) \rho(t,y) \overset{J_2}{(u(t,x) - u(t,y))} \partial_t (u(t,x) - u(t,y)) dy$$

Then we use the equations satisfied by the macroscopic equations

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$$\begin{cases} \partial_t \rho + \operatorname{div}_x(\rho u) = 0 \\ \rho \partial_t u + \rho u \cdot \nabla_x u + \operatorname{div}_x P = \int L(\beta) f dv \end{cases}$$

and get that

$$\begin{aligned} J_1 &= 4 \int \rho u \cdot \nabla_x u \cdot u \, dx \\ &\quad - 4 \left( \int \rho u \cdot \nabla_x u \, dx \right) \cdot \int \rho u \, dx \end{aligned}$$

$$\begin{aligned} J_2 &= -J_1 - 4 \int \operatorname{div} P \cdot u \, dx \\ &\quad - 2 \int \Psi(x-y) \rho(t,x) \rho(t,y) |u(t,x) - u(t,y)|^2 \, dx dy \end{aligned}$$

Therefore we have

$$\frac{dE_2}{dt} = -4 \int \operatorname{div} P \cdot u - 2 \int \Psi(x-y) \rho(t,x) \rho(t,y) |u(t,x) - u(t,y)|^2 \, dx dy$$

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Gathering the two results, it yields

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$$\begin{aligned}\frac{d\varepsilon}{dt} &\leq -2\varepsilon_1 - \Psi_m \varepsilon_2 \\ &\leq -2 \min(1, \Psi_m) \varepsilon\end{aligned}$$

we conclude by a Gronwall lemma  $\square$