

# Spectral stability and singular perturbation

Louis Garénaux

*Institut de Mathématiques de Toulouse*

Montie Avery

*University of Minnesota*

Séminaire EDP, IRMAR

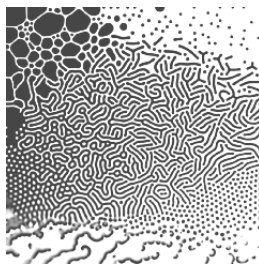
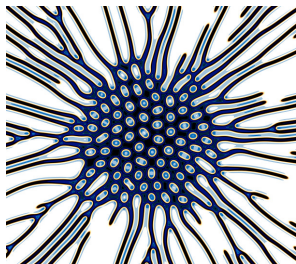
07/01/21

# Reaction-difusion equation

$$\partial_t u = Du + R(u)$$

Examples: Gray-Scott equations:

$$\begin{cases} \partial_t u = \Delta u - uv^2 + a(1 - u), \\ \partial_t v = d\Delta v + uv^2 + bv, \end{cases}$$



<https://www.karlsims.com/rd.html>

Fitzhugh-Nagumo, Allen-Cahn, bistable equations

# A simple equation, and a perturbation

- The scalar, one dimensional FKPP equation writes:

$$\partial_t u = \partial_{xx} u + u(1 - u), \quad (\text{FKPP})$$

where  $u(t, x) \in \mathbb{R}$ ,  $t > 0$  and  $x \in \mathbb{R}$ . First studied in 1937, it models the spreading of a population (bacteria, muskrat, chemical reaction).

- The extended FKPP equation writes:

$$\partial_t u = -\delta^2 \partial_{xxxx} u + \partial_{xx} u + u(1 - u). \quad (\text{eFKPP})$$

It can be obtained as an amplitude equation for the solutions of a reaction-diffusion system, that undergoes a (co-dimension 2) bifurcation. In our case, it is a good model to see how properties depend on the equation.

For both problem, there exists front-like solutions that travel at constant speed  $c$ , greater than the critical speed  $c_* = 2 + O(\delta)$ . See [KPP 1937](#) and [Rottschäfer–Wayne 2001](#).





# Asymptotic stability of the front

We see the critical front  $q_*$  as an equilibrium, and study its stability with respect to perturbations  $p \in H^4(\mathbb{R})$ .

There exists  $\varepsilon > 0$  such that for any  $p_0 \in X$  that satisfy  $\|p_0\| \leq \varepsilon$ , the solution  $u(t, x) = q_*(x) + p(t, x)$  is defined for all time, and the perturbation satisfy

$$\|p(t, \cdot)\| \leq f(t)\|p_0\|,$$

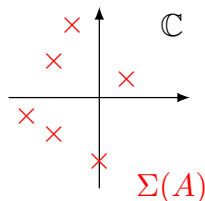
where  $f(t) \rightarrow 0$  when  $t \rightarrow +\infty$ .

- 1976 Sattinger: the supercritical fronts for FKPP are exponentially stable in weighted spaces. 
- 1994 Gallay: the critical front for FKPP is algebraically stable in a weighted space:  $f(t) = \frac{1}{(1+t)^{3/2}}$ . Furthermore, the decay is optimal. 
- 2018 Faye–Holzer: re-obtain the algebraic decay of Gallay by more robust method. The weak decay of the front  $q'_*(x) \sim xe^{-x}$  at  $+\infty$  is essential.
- 2001 Rottschäfer–Wayne: The critical front for eFKPP is stable in weighted space. No rate of convergence.

# A dynamical point of view

Let  $F : \mathbb{R}^n \rightarrow \mathbb{R}^n$  and  $\bar{y} \in \mathbb{R}^n$  such that  $F(\bar{y}) = 0$ . The dynamic near the equilibrium is given by the spectrum of the matrix  $A = dF(\bar{y})$ . In our case, we want to study

$$A(\delta) = -\delta^2 \partial_{xxxx} + \partial_{xx} + c_*(\delta) \partial_x + (1 - 2q_*(\delta, x)).$$



## Theorem (Avery–Scheel 2020)

If the linear dynamic satisfy the spectral hypotheses 1 through 4, then for any initial condition  $p_0$  with  $\|\omega p_0\|_{H_r^1(\mathbb{R})}$  small enough, the solution  $u(t, x) = q_*(x) + p(t, x)$  of the nonlinear problem satisfy

$$\|\omega p(t, \cdot)\|_{H_{-r}^1} \leq \frac{C}{(1+t)^{3/2}} \|\omega p_0\|_{H_r^1}.$$

Furthermore the rate is optimal. There exists  $\psi$  that does not depends on  $p_0$  such that  $\|\omega p(t, \cdot) - t^{-3/2} \psi\|_{H_{-r}^1} \leq \frac{C}{(1+t)^2} \|\omega p_0\|_{H_r^1}$ .

# Definitions

We want to study the spectrum of

$$\mathcal{A}(\delta) = -\delta^2 \partial_{xxxx} + \partial_{xx} + c_*(\delta) \partial_x + (1 - 2q_*(\delta, x)),$$

which is defined on the whole real line:  $\mathcal{A}(\delta) : H^4(\mathbb{R}) \subset L^2(\mathbb{R}) \longrightarrow L^2(\mathbb{R})$ .

$$\Sigma(\mathcal{T}) = \{\lambda \in \mathbb{C} : \lambda - \mathcal{T} \text{ is not bounded invertible}\}.$$

It is the disjoint union of the point spectrum  $\Sigma_{\text{pt}}(\mathcal{T})$ :

$$\dim \ker(\lambda - \mathcal{T}) - \text{codim im}(\lambda - \mathcal{T}) = 0,$$

and the essential spectrum  $\Sigma_{\text{ess}}(\mathcal{T})$ :

$$\dim \ker(\lambda - \mathcal{T}) - \text{codim im}(\lambda - \mathcal{T}) \neq 0.$$

Furthermore, we say that  $\lambda$  is:

- an eigenvalue if there exists  $u \in H^4(\mathbb{R})$  such that  $\mathcal{T}u = \lambda u$ ,
- a resonance if there exists  $u \in W^{4,\infty}(\mathbb{R})$  such that  $\mathcal{T}u = \lambda u$ .

# Examples

- Differential operator with constant coefficients:

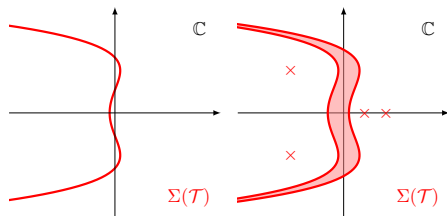
$$\mathcal{T} = a_m \partial_x^m + \cdots + a_1 \partial_x + a_0 = P(\partial_x)$$

The essential spectrum is the curve  $P(i\mathbb{R}) = \{P(ik) : k \in \mathbb{R}\}$ . There is no point spectrum.

- Coefficients that converge at exponential speed to distinct limits at  $\pm\infty$ :

$$\mathcal{T}(x) = a_m(x) \partial_x^m + \cdots + a_1(x) \partial_x + a_0(x),$$

where both  $\sup_{x>0} |a_i(x) - a_i^+| e^{\alpha|x|}$  and  $\sup_{x<0} |a_i(x) - a_i^-| e^{\alpha|x|}$  are finite. Then the essential spectrum is contained in between the two curves  $\Sigma(\mathcal{T}^+)$  and  $\Sigma(\mathcal{T}^-)$ . The point spectrum is not (necessarily) empty.



## Sturm–Liouville theory

If  $m = 2$ , and if  $a_i(x)$  are real, then  $\Sigma_{\text{pt}}(\mathcal{T})$  consists of a finite number of *simple* eigenvalues  $\lambda_0 > \cdots > \lambda_N$ , and the associated eigenfunction  $u_i$  vanishes exactly  $i$ -times.

# Main Theorem

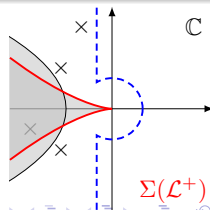
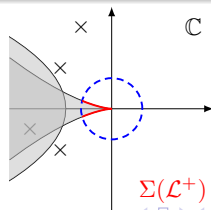
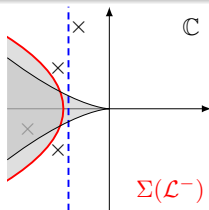
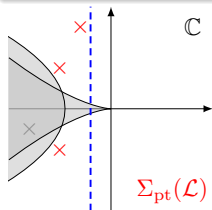
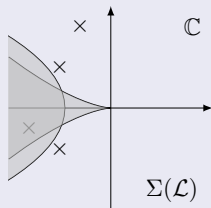
## Theorem (Avery–G. 2020)

There exists  $\delta_0 > 0$  such that for all  $0 < \delta < \delta_0$ , the linear dynamic of (eFKPP) in a weighted space  $\mathcal{L} = \omega \mathcal{A} \omega^{-1}$ , satisfies the following spectral hypotheses.

1. Stable point spectrum and stable spectrum for  $\mathcal{L}^-$ .
2. Simple pinched double root at  $(\lambda, k) = (0, 0)$  for  $\mathcal{L}^+$ :

$$P^+(ik) = -\alpha k^2 + O(k^3).$$

3. Stable spectrum for  $\mathcal{L}^+$ , away from the origin.
4. No eigenvalue and no resonance at  $\lambda = 0$ .





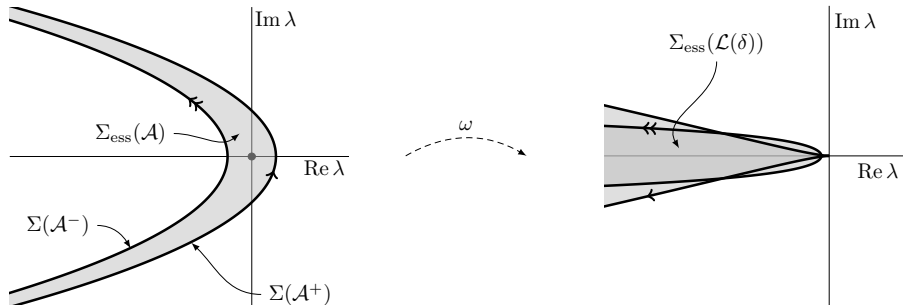
# Essential spectrum for eFKPP

## Proposition

There exists a weight  $\omega$  such that  $\mathcal{L} := \omega \mathcal{A} \omega^{-1}$  has marginally stable essential spectrum.

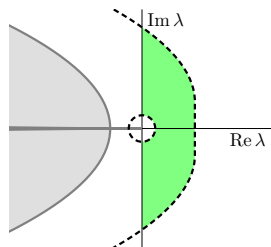
*Proof:* Choose  $\omega$  of the form  $\omega(x) = \begin{cases} e^{\eta_* x} & x \geq 1, \\ 1 & x \leq -1. \end{cases}$

Then the spectrum at  $-\infty$  is unchanged:  $\Sigma(\mathcal{L}^-) = \Sigma(\mathcal{A}^-) = P^-(i\mathbb{R})$ , while the spectrum of  $\mathcal{L}^+$  is the curve  $P^+(-\eta_* + i\mathbb{R})$ .  $\square$



# Point spectrum for eFKPP – step 1

We decompose the study in three parts. Near the origin, outside of a compact set, and finally the intermediate domain.



## Proposition

There exists a compact  $K \subset \mathbb{C}$  such that  $\mathcal{L}(\delta)$  has no unstable eigenvalues outside of  $K$ .

*Proof:* Use the ellipticity of  $\mathcal{L}(\delta)$ : assume that  $\mathcal{L}(\delta)u = \lambda u$  and take the scalar product with  $u$ . It leads to estimates on  $\operatorname{Re} \lambda$  and  $\operatorname{Im} \lambda$ . □

# Point spectrum for eFKPP – step 2

In the intermediate domain (green part), we treat  $\mathcal{L}(\delta)$  as a perturbation of  $\mathcal{L}(0)$ . From Sturm-Liouville theory, the second order operator  $\mathcal{L}(0)$  has no unstable point spectrum.

## Proposition

There exists  $\delta_0 > 0$  and  $r(\delta_0) > 0$  such that if  $0 < \delta < \delta_0$ , then the operator  $\mathcal{L}(\delta)$  has no eigenvalue inside  $K \setminus B(0, r)$ .

*Proof:* Remove the singularity by applying the preconditioner  $(1 - \delta^2 \partial_{xx})^{-1}$ . Computations leads to

$$(1 - \delta^2 \partial_{xx})^{-1}(\mathcal{L}(\delta) - \lambda)u = (\mathcal{L}(0) - \lambda)u + e(\delta, u),$$

where the error term  $e(\delta, u)$  goes to 0 in  $L^2(\mathbb{R})$  as  $\delta \rightarrow 0$ .

Since  $\lambda$  is away from the spectrum of  $\mathcal{L}(0)$ , we can invert  $\mathcal{L}(0) - \lambda$ . □

# Point spectrum for eFKPP – step 3

When perturbing an operator, point spectrum can form near the essential spectrum. The previous step can not apply at the origin. Instead, we use a Lyapounov-Schmidt decomposition, as done in [Pogan-Scheel 2010](#).

## Proposition

There exists an holomorphic function  $E : B(0, \delta_0) \times B(0, r_0) \rightarrow \mathbb{C}$ , such that the eigenproblem

$$\mathcal{L}(\delta)u = \lambda u$$







admits a bounded solution  $u$  if and only if  $E(\delta, \sqrt{\lambda}) = 0$ . Furthermore,  $E(0, 0) \neq 0$ . In particular, there is no resonance at the origin.

*Proof:* Use implicit function theorem on a suitable formulation of the eigenproblem. Noting  $v := (1 - \delta^2 \partial_x^2)^{-1} (\mathcal{L}(\delta) - \lambda)u$ , the problem  $v = 0$  is equivalent to

$$P_{\text{im } \mathcal{L}(0)} v = 0, \quad \text{and} \quad \langle v, \varphi \rangle = 0,$$

due to Fredholm property of  $\mathcal{L}(0)$ , and strong localization of  $u$ . We can solve the first equation locally, and use the second equation to define  $E(\delta, \sqrt{\lambda})$ .  $\square$

## Thank's for your attention!

-  M. Avery and L. Garénaux. Spectral stability of the critical front in the extended fisher-kpp equation, 2020.
-  M. Avery and A. Scheel. Asymptotic stability of critical pulled fronts via resolvent expansions near the essential spectrum, 2020.
-  G. Faye and M. Holzer. Asymptotic stability of the critical fisher–kpp front using pointwise estimates. *Zeitschrift für angewandte Mathematik und Physik*, 70(1), Nov 2018.
-  T. Gallay. Local stability of critical fronts in nonlinear parabolic partial differential equations. *Nonlinearity*, 1994.
-  A. Pogan and A. Scheel. Instability of spikes in the presence of conservation laws. *Zeitschrift für angewandte Mathematik und Physik*, 61:979–998, 12 2010.
-  V. Rottschäfer and C. E. Wayne. Existence and stability of traveling fronts in the extended Fisher-Kolmogorov equation. *Journal of Differential Equations*, 2001.